EXPERIMENTAL AND COMPUTATIONAL STUDIES ON THE TEMPERATURE DEPENDENCE OF THERMAL CONDUCTIVITIES FOR WATER, ETHYLENE GLYCOL, GLYCEROL, AND PROPYLENE GLYCOL, USING THE TRANSIENT HOT-WIRE METHOD

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Authorization to Submit Thesis

This thesis of Salman M. Alharbi, submitted for the degree of Master of Science with a Major in Mechanical Engineering and titled "Experimental and Computational Studies on the Temperature Dependence of Thermal Conductivities for Water, Ethylene Glycol, Glycerol, and Propylene Glycol, Using the Transient Hot-Wire Method," has been reviewed in final form. Permission, as indicated by the signatures and dates below, is now granted to submit final copies to the College of Graduate Studies for approval.

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Abstract

The goal of this work is to provide experimental measurements of thermal conductivity of water, ethylene glycol, glycerol, and propylene glycol as a function of temperature. The transient hot wire method was used to measure the thermal conductivity over temperatures ranging from 235–340 K. This work also involved in-house apparatus fabrication along with integration of data acquisition and processing software. The experiments are carried out for a fixed current of 250 mA and the resulting temperature rise of a 95.33 mm long, 25-micron radius platinum wire is used to infer the thermal conductivity using the known solution to the heat conduction equation for a continuous line source in an infinite medium. It is important to account for the variable temperature coefficient of resistance of the platinum wire as a function of temperature when seeking to obtain the correct temperature dependence of the thermal conductivity. A data reduction procedure that improves the accuracy of the reported values by identifying the onset of convection in the fluid is proposed. We use the peak value of the slope (S) obtained using a third order polynomial fit to the apparent linear region to estimate the thermal conductivity. The high-resolution data acquired at closely spaced temperature intervals is used to derive a correlation between thermal conductivity values and the fluid temperature. Additionally, numerical results for temperature and velocity field near the heated wire are also presented to help understand the non-idealities present in the experiments. The experimental temperature rise obtained from the transient hot-wire experiments is compared to computed values for water at room temperature, and a good agreement is found. There is a fair agreement between the current data sets and the very limited data for the four liquids reported in the literature. This work provides robust and comprehensive experimental data for thermal conductivities of the four common heat transfer fluids over the typical range of temperatures they are frequently used.

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Dedication

Dedicated to the memory of my mother.

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Table of Contents

Authorization to Submit Thesis	ii
Abstract	iii
Acknowledgements	iv
Dedication	v
List of Tables	viii
List of Figures	ix
Chapter 1. Introduction	
1.1 Heat Transfer—A Historical Perspective	
1.2 Conduction Heat Transfer	
1.3 Thermal Conductivity	7
1.4 Measurement Techniques for Thermal Conductivity	
1.4.1 Steady State Methods	
1.4.2 Transient Methods	9
1.5 Objectives of this Work	9
Chapter 2: The Transient Hot-Wire Technique for Measuring Thermal Conductivity .	
2.1 Introduction and Theory	
2.2 Apparatus	
2.3 Determination of Temperature Coefficient of Resistance	
2.4 Data Reduction Procedure	
Chapter 3: Temperature Dependence of Thermal Conductivity for Water: Experiment	ts and
2.1 Thermal Conductivity Deputts for Water	
2.1.1 Comparison with Providuals Perpented Data	
2.2.C to the the	
3.2 Computed Results	
3.2.1 Computational Specifications and Results	
3.2.2 Temperature Distribution and Local Heat Transfer Coefficients	
3.2.3 Velocity Field	
3.3 Comparison of Experimental and Computed Results	
3.4 Summary	

Chapter 4: Temperature Dependence of Thermal Conductivity for Ethylene Glycol and Propylene Glycol	, Glycerol,
4.1. Ethylene Glycol	
4.1.1 Thermal Conductivity Results for Ethylene Glycol	
4. 2 Glycerol	41
4.2.1 Thermal Conductivity Results for Glycerol	44
4.3 Propylene Glycol	
4.3.1 Thermal Conductivity of Propylene Glycol	49
4.4 Comparative Thermal Conductivities for the Four Liquids	52
Chapter 5: Conclusions and Future Work	55
5.1 Summary of this Study	55
5.2 Future Work	56
References Cited	
Appendix A: Tabulated Thermal Conductivity Results for Water	
Appendix B: Tabulated Thermal Conductivity Results for Ethylene Glycol	66
Appendix C: Tabulated Thermal Conductivity Results for Glycerol	
Appendix D: Tabulated Thermal Conductivity Results for Propylene Glycol	

List of Tables

Table 1-1 Contributions leading to the heat diffusion equation [11]	. 6
Table 3-1 Material Properties used in the simulations	27

List of Figures

Figure 1.1: Heat Flow in a one-dimensional rod
Figure 1.2: Variation of thermal conductivity with temperature for a solid [12], liquid [13], and gas [14]7
Figure 2.1: Schematic of the transient hot-wire cell
Figure 2.2: Variation of platinum wire resistance with temperature
Figure 2.3: Determination of local Slope
Figure 2.4: Variation of the temperature coefficient of resistance for platinum wire with temperature
Figure 2.5: Temperature-time history for experimental runs at varying liquid temperatures 19
Figure 2.6: Illustration of the three distinct temperature rise regions observed in the experiments
Figure 2.7: (a) Region of experiment used for fitting and extracting the slope (b) Comparison of slopes obtained using a linear and a 3rd order polynomial fit. The filled symbol represents the slope used in the experimental determination of thermal conductivity
Figure 3.1: Three sets of experimental runs showing thermal conductivity variation with temperature for water
Figure 3.2: Consolidated data set for thermal conductivity as a function of temperature for current experimental sets for water
Figure 3.3: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data
Figure 3.4: Comparison of current data with other recent studies for water
Figure 3.5: (a) A zoomed in view of the grid distribution. Note that the total simulated domain extends to 150 and 600 non-dimensional units in the axial (x/a) and radial (r/a) directions respectively. Here <i>a</i> is the radius of the wire and <i>g</i> the acceleration due to gravity. (b) Streamtraces and temperature contours at t=6.0 s
Figure 3.6: Computed contours of non-dimensional (a) Temperature rise at $t = 6s$, and (b) the radial temperature distribution midway at $x/a = 75$ at various times
Figure 3.7: (a) Scaled local heat transfer coefficient (a) as a function of distance (b) Fit (solid line) to the steady state distribution as a function of distance, x' , from the entrance
Figure 3.8: Scaled local heat transfer coefficient at $xa = 75$ as a function of time

Figure 3.9: Non-dimensional radial temperature and axial velocity profiles at $x = L/$	2 33
--	------

Figure 3.10: Comparative temperature-time history for (a) Experiments versus CFD results, and (b) Exact solution to heat diffusion equation for a line source in an infinite media versus CFD results
Figure 4.1: Three sets of experimental runs showing thermal conductivity variation with temperature for ethylene glycol
Figure 4.2: Consolidated data set for thermal conductivity as a function of temperature for current experiments for ethylene glycol
Figure 4.3: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for ethylene glycol
Figure 4.4: Comparison of current data with other recent studies for ethylene glycol
Figure 4.5: Three sets of experimental runs showing thermal conductivity variation with temperature for glycerol
Figure 4.6: Consolidated data set for thermal conductivity as a function of temperature for current experimental sets for glycerol
Figure 4.7: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for glycerol
Figure 4.8: Comparison of current data with other recent studies for glycerol
Figure 4.9: Three sets of experimental runs showing thermal conductivity variation with temperature for propylene glycol
Figure 4.10: Consolidated data set for thermal conductivity as a function of temperature for current experimental sets for propylene glycol
Figure 4.11: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for propylene glycol
Figure 4.12: Comparison of current data with other recent studies for propylene glycol 52
Figure 4.13: Comparative thermal conductivities of all four liquids obtained in this work 53
Figure 4.14: Plot showing the variability in the thermal conductivity values for the four liquids in the neighborhood of 298 ± 1 K

Chapter 1. Introduction

1.1 Heat Transfer—A Historical Perspective

The flow of energy because of temperature differences forms the basis for the discipline of heat transfer. The distinction between heat and temperature was not clear during the 18th century. An anonymously published work titled "Scala Graduum Caloris" [1] was an initial attempt to establish a temperature scale. According to Sayre [2], the title of this article first published in 1701 had been translated as "Scales of the degree of heat", and uses the present concepts of 'heat' and 'temperature' interchangeably. There is some evidence to claim that the author of this work was probably Sir Isaac Newton. The scale for the degree of calor in the low range was determined using the principles of thermal expansion of liquid, with the melting calor of lead being the upper limit. The process for determining the higher degrees of calor involved the measurement of time required cooling of various combinations of different metal pieces placed on an initially glowing iron bar. The mention of the second method in this work, in the opinion of Sayre [2], is the only link between 'the law of cooling' and Newton. A key postulate of this work was that the uniform flow of air removed calor from the hot body in proportion to the calor difference.

In the later part of the eighteenth century, two rival theories of heat emerged. The first was the caloric theory of heat proposed by Antoine Lavoisier [3]. The concept of the caloric theory was that the heat was an invisible subtle fluid whose particles were in motion [4]. The fluid was tasteless, odorless, massless, and colorless. This fluid was transferred from a hot to a cold body during the heating process. This idea was popular among the chemists and presupposed the existence of atoms surrounded by the fluid. The caloric theory was successfully used to explain all known heat related phenomena and was accepted by scientists such as

Laplace, Lavoisier, Priestley, Petit and Dulong [5]. Based on this (incorrect) theory the French engineer Sadi Carnot (correctly) deduced the fundamental limitations of conversion of heat to work [6]. The second theory, accepted among physicists and mathematicians, stated that there was no such fluid and that the motion of 'atoms' could account for heat. Interestingly, the first law of thermodynamics and the atomic theory of matter were still unknown at the time. The cannon boring experiments of Count Rumford provided conclusive evidence against the caloric theory. A blunt cannon borer was seemingly able to provide a limitless amount of the caloric fluid. He stated the problem as [5]:

"Whence then came this heat? And what is heat actually? I must confess that it has always been impossible for me to explain the results of such experiments except by taking refuge in the very old doctrine which rests upon the supposition that heat is nothing but a vibratory motion taking place among the particles of the body".

This idea helped to develop the law of conservation of energy, and J.P Joule established the equivalence between mechanical work and heat in 1842 [7]. This equivalence is what is now referred to as the First Law of Thermodynamics. Two decades prior, Fourier's work on the Analytic Theory of Heat outlined the basic principles of heat transfer [8]. His work showed that the flow of heat was a result of differences in temperature between adjacent particles, or simply put due to spatial temperature gradients. The principles of conduction heat transfer had now been firmly established.

1.2 Conduction Heat Transfer

Heat transfer is thermal energy in motion. It involves exchange in thermal energy from one object to another object because of temperature differences. Thermal energy is always transferred from a high temperature to the low temperature object until both reach the same temperature. The origin of conduction heat transfer can be traced to molecular vibrations in a medium where there is no bulk motion. This mode can exist in either solids, liquids, or gases. For the case of solids the transfer of heat is attributed to lattice vibrations and electronic contributions, whereas for liquids and gases the random molecular motions are important [9].

For simplicity, let us consider the simple example of heat flow due to conduction in a solid onedimensional rod. Following the analysis in the text of Haberman [10] let us consider a rod of length L aligned with the x-axis. Let the thermal energy density be given by:

$$e(x,t) \equiv$$
 thermal energy density 1.1

Let the flow of thermal energy, i.e., the quantity of thermal energy flowing per unit time per unit area equal to

$$\phi(x,t) = \text{Heat Flux}$$

$$A \quad \phi(x,t) \qquad \phi(x + \Delta x, t) \qquad \\ x = 0 \qquad x + \Delta x \qquad x = L$$

Figure 1.1: Heat Flow in a one-dimensional rod.

For generality, let us assume that the rod possesses some source of internal energy generation whose strength (heat generated per unit volume per unit time) is given by

1.2

$$Q(x, t)$$
 = Heat generated per unit volume per unit time 1.3

The conservation of energy principle applied to the slice of length Δx can be stated as:

$$\begin{pmatrix} \text{Rate of change} \\ \text{of heat energy} \\ \text{in time} \end{pmatrix} = \begin{pmatrix} \text{heat energy flowing} \\ \text{across boundaries} \\ \text{per unit time} \end{pmatrix} + \begin{pmatrix} \text{Internal generation of} \\ \text{heat energy} \\ \text{per unit time} \end{pmatrix}$$
 1.4

For a thin slice of width Δx the total thermal energy generated is given by $Q(x,t)A\Delta x$ while the heat energy contained within it is $e(x,t)A\Delta x$. Here we are assuming negligible variation in e(x,t) and Q(x,t) over the thin slice. The conservation principle stated in equation 1.4 can be more precisely stated as

$$\frac{\partial(e(x,t)A\Delta x)}{\partial t} = \phi(x,t)A - \phi(x+\Delta x)A + Q(x,t)A\Delta x$$
 1.5

Dividing throughout by Δx and taking the limit $\Delta x \rightarrow 0$, we obtain

$$\frac{\partial e}{\partial t} = \lim_{\Delta x \to 0} \frac{\phi(x,t) - \phi(x + \Delta x,t)}{\Delta x} + Q(x,t)$$
 1.6

While taking the limit $\Delta x \rightarrow 0$ the time is held constant. Therefore, by the definition of the partial derivative one can simplify the above equation to

$$\frac{\partial e}{\partial t} = -\frac{\partial \phi}{\partial x} + Q \tag{1.7}$$

Equation 1.7 contains two variables e (thermal energy density) and ϕ (heat flux) which must be represented or modeled in terms of some common variable. Physically, this variable is called the temperature T. We must therefore provide a relationship between the temperature and the thermal energy. This relationship can be shown to be of the form

$$e(x,t) = \rho(x)c(x)T(x,t)$$

Where ρ and *c* represent the density and specific heat. This relationship between thermal energy and temperature had confounded the early practitioners of the science of heat, and the notions of heat and temperature were often used interchangeably [2]

Next, we need to model the heat flux ϕ in terms of the temperature field. This will reduce equation 1.7 to one unknown, namely *T*. Fourier established the relationship between the temperature field and the heat flux. He initially formulated heat conduction as an *n*-body problem similar to Biot who had worked on the problem it before him. Biot, who belonged to the Laplace's school of thought, adhered to the principle of action at a distance between bodies. His concept involved only the temperature difference between points, and that the temperature at a given point was influenced by all neighboring points. Critically, it did not involve distances and hence the temperature gradients. Fourier, subsequently abandoned the discontinuous *n*-body approach and adopted a continuous approach. He assumed that the temperature within an infinitesimal element was affected only by elements in its immediate vicinity, and consequently formulated the heat diffusion equation for a continuum. His empirical approach involved spatial transport of heat, storage of heat, and the interaction of the domain with the boundary conditions [11]. This approach also led to what is known as Fourier's law of heat conduction which relates the heat flux ϕ to the spatial temperature gradient as

$$\vec{\phi} = -k\vec{\nabla}T \tag{1.8}$$

The quantity k is proportionality constant called the thermal conductivity and is a material dependent property. Fourier's 1807 work on the heat diffusion equation faced delays in publication and he would finally publish it himself as Théorie Analytique de la Chaleur in 1822. The delay resulted from his use of infinite trigonometric series in the solution of the problem which was not viewed favorably by Lagrange, one of the reviewers of his 1807

manuscript submitted to the French Academy [11]. The transient heat conduction equation of Fourier in the absence of heat sources is given by:

$$\rho c \frac{\partial T}{\partial t} = \vec{\nabla} \cdot k \vec{\nabla} T \tag{1.9}$$

	Year	Contribution
Fahrenheit	1724	Mercury thermometer and standardized temperature scale
Abbé Nollet	1752	Observation of osmosis across an animal membrane
Bernoulli	1752	Use of trigonometric series for solving differential equations
Black	1760	Recognition of latent heat and specific heat
Crawford	1779	Correlation between respiration of animals and their body heat
Lavoisier and Laplace	1783	First calorimeter; measurement of heat capacity, latent heat
Laplace	1789	Formulation of Laplace operator
Biot	1804	Heat conduction among discontinuous bodies
Fourier	1807	Partial differential equation for heat conduction in solids
Fourier	1882	Théorie Analytique de la Chaleur

Table 1-1 Contributions leading to the heat diffusion equation [11].

As noted earlier the heat diffusion equation includes the thermal conductivity term, which appears as a consequence of modeling the heat flux term as a function of the temperature field (gradient). Formally, the definition of thermal conductivity is based on Fourier's law of conduction and for a one-dimensional case can be written as:

$$k \equiv -\frac{\phi_x}{\partial T / \partial x}$$
 1.10

Where ϕ_x is the heat flux and $\partial T/\partial x$ is the temperature gradient. In order to utilize Fourier's law of conduction, we need to have knowledge of the thermal conductivity of the material. This material property may vary as a function of position for anisotropic solids, and

is often a function of temperature. It fundamentally relates to the idea of transport of heat and is a measure of the rate at which heat is transferred by diffusion [9].

1.3 Thermal Conductivity

The transport of heat in the solid-state occurs due to contributions arising from two distinct processes. First the movement of free electrons and the other due to lattice vibration waves. These two effects are additive in that the observed thermal conductivity is the sum of these two components. Thermal conductivity for metallic solids generally tends to be much higher than nonmetals where the thermal conductivity arises primarily due to the lattice vibrations. In the case of liquids and gases theory states that the thermal conductivity bears a direct proportionality to factors such as particles per unit volume (*n*), the mean molecular speed (\bar{c}), and the mean free path (λ), i.e. $k \propto n\bar{c}\lambda$ [9]. The effectiveness of thermal transport can vary over several orders of magnitude between solids, liquids, and gases as shown in Fig. 1.2.



Figure 1.2: Variation of thermal conductivity with temperature for a solid [12], liquid [13], and gas [14].

Furthermore, the thermal conductivity exhibits a strong temperature dependence, which may cause it to either increase or decrease with increasing temperature depending on the material.

It is therefore of fundamental and practical interest to obtain the temperature dependence of this transport property in order to optimally design engineering systems such as heat exchangers that typically operate over a wide temperature range.

1.4 Measurement Techniques for Thermal Conductivity

There are several experimental techniques to measure thermal conductivity of materials. All of these methods can be fundamentally characterized as either steady state or transient. In this work, we use one of the transient techniques known as the transient-hot wire method to obtain the thermal conductivity as a function of temperature. A brief description of the two classes of techniques is provided in the following.

1.4.1 Steady State Methods

The steady state methods mean that the thermal conductivity of sensing device does not change with time. A commonly used steady state method is known as the concentric cylinder method. In this method the gap between the two cylinders is filled with the test material and heat is supplied to the inner cylinder and flows across the specimen in the gap. Thermocouples placed located at the inner surface of the outer cylinder r_2 , and the outer surface of the inner cylinder r_1 , record the temperature difference across the gap. Under steady state conditions and for a known heat supply Q_n the thermal conductivity is determined as:

$$k = \frac{\ln\left(\frac{r_2}{r_1}\right) \times Q_n}{2\pi L \Delta T}$$
 1.11

Here L is the length of the inner cylinder; ΔT is the different temperatures. The net power Q_n is equal to the total power supplied minus the heat loss. In general, the steady state

methods requires a long test time because the temperature one must let the system to first achieve steady state.

1.4.2 Transient Methods

The most commonly used transient technique in the literature is the transient hot-wire apparatus. Other techniques such as the transient plane source and the laser flash method also are used to determine the thermal conductivity of substances. The fundamental idea in these measurements is to quickly deposit a short pulse of energy in the medium and record the transient temperature response. This temperature-time response is then compared to fundamental solutions and a curve fit procedure is often used to estimate the thermal conductivity. Oftentimes, the heat source and the sensing material are the same such as a thin platinum wire. We will discuss the transient hot-wire technique in detail in Chapter 2.

1.5 Objectives of this Work

- Design and construct a transient hot-wire apparatus for the measurement of thermal conductivity of liquids. The design of the overall measurement system involves construction of a hot-wire cell, integration of heating and cooling baths to the system with precision temperature control, a precision power source, and a high-resolutionhigh-speed data acquisition device. In addition to the hardware itself, a key component of the apparatus is the software that enables high-speed data acquisition and processing in near real-time.
- 2. Obtain the temperature dependence of thermal conductivity of four common heat transfer fluids, namely, water, ethylene glycol, glycerol, and propylene glycol. The thermal conductivity values obtained from the system were validated using water, which is the most widely available heat transfer fluid and has extensively been studied with

respect to the variation of its thermal conductivity with temperature. Post validation, we obtain results for the other three liquids that do not have a large volume of work on their thermal conductivity values. In this work, we wish to generate sufficient number of data points for thermal conductivity over the temperature range of interest so that a robust correlation of k vs. T can be obtained.

3. Computationally examine the temperature and velocity field in the vicinity of the transient hot-wire. The simulations are required to understand the physics of the heat up process for the micron sized wire and its interaction with the surrounding fluid media as it pertains to the onset of buoyancy driven flow. The computational results provide insights into the validity of assumptions made in the process of measuring thermal conductivity of liquids using this method.

Chapter 2: The Transient Hot-Wire Technique for Measuring Thermal Conductivity

2.1 Introduction and Theory

Thermal conductivity of a heat transfer liquid is among the important properties required for proper estimation of heat transfer rates in practical engineering systems. This fundamental transport property relates the temperature gradient to heat flux. A widely used technique used for the measurement of thermal conductivity is the transient hot-wire method. This method is based on the solution to the conduction problem $(\partial T/\partial t = \kappa \nabla^2 T)$ of radial heat flow in an infinite solid medium with an instantaneous line source provided with a constant supply of heat [15]. In practice, a very thin platinum wire supplied with a constant current, and immersed in a liquid adequately approximates the ideal configuration. Assuming that the power is applied to line source at t = 0, and that the wire and the surrounding medium at the same temperature, the temperature at any subsequent time t and at a distance r is given by:

$$\Delta T(r,t) = \frac{q}{4\pi k} \int_{r^2/4\kappa t}^{\infty} \frac{e^{-u}}{u} du$$
(2.1)

Of particular interest is the temperature at the wire surface (r = a). The transient hot-wire method for measuring thermal conductivity relies on the fact that at large time $(Fo = \frac{\kappa t}{a^2} \gg 1)$ the relationship between the temperature rise at the wire surface and the logarithm of time is linear. The slope of this linear fit is inversely related to the thermal conductivity for a constant heating rate. Specifically, the temperature rise of the wire (r = a) and for $t \gg a^2/\kappa$ can be approximated as [16]:

$$\Delta T(a,t) = \frac{q}{4\pi k} ln\left(\frac{4\kappa t}{a^2 C}\right)$$
(2.2)

Here a is the radius of the wire, q the power per unit length supplied to the wire, κ and k the thermal diffusivity and conductivity of the medium surrounding the wire, and $C = e^{\gamma}$ where $\gamma = 0.577216$ is Euler's constant. The derivative of the temperature rise with respect to logarithm of time leads to:

$$\frac{d\Delta T}{dln(t)} = \frac{q}{4\pi k} \tag{2.3}$$

Referring to the quantity $d\Delta T/dln(t)$ as the slope S of the plot of ΔT versus ln(t) the thermal conductivity of the surrounding medium is inferred using:

$$k = \frac{q}{4\pi S} \tag{2.4}$$

Additional details on the theory of the transient hot-wire method can be found in the studies by Blackwell [17], Healy et al.[18], and Jaeger [15]

The application of this technique in experiments for determining the thermal conductivity of fluids has been described by Roder [19] who designed a new apparatus for measuring thermal conductivity of oxygen at elevated pressures. He also performed performance checks using nitrogen, helium and argon. His work also provides a brief overview of the evolution of this technique starting from the initial experiments of Pittman [20]. A method for simultaneous measurement of thermal diffusivity and conductivity for liquids has been reported by Nagasaka and Nagashima [21] who carried out measurements on toluene under atmospheric pressure and in the temperature range of 0 - 80 C. Among the recent studies describing the use of the transient hot-wire apparatus for liquid thermal conductivity measurements are the studies by Bleazard et al. [22], Codreanu et al. [23], Zhang et al. [24], and Kostic and Simham [25].

The accurate experimental determination of thermal conductivity has generally been regarded as a problem of some difficulty [26]. One of the factors that tends to overwhelm

experimental efforts of measurements in fluids is the onset of convection. Note that the aforementioned solution holds for heat diffusion in a solid/stationary-fluid medium, and it breaks down if natural convection sets in due to heat input to the system. This work reports the thermal conductivity of water in the 273 - 305 K temperature range at a pressure of one atmosphere. Prior work on the temperature dependence of thermal conductivity of water includes the study of Woolf et al. [27] who used a concentric cylinder apparatus under steady state conditions and between 70 and 200 Fahrenheit. Lawson et al.[28] reported the thermal conductivity of water between 30 and 140 degree Celsius and for pressures between 1 to 8000 kg/cm2, also using the concentric cylinder technique. Theiss and Thodos [29] reviewed the experimental thermal conductivity and viscosity for gaseous and liquid water and developed a reduced state correlation for these transport properties. They note that while there is significant experimental work for viscosity, the thermal conductivity measurements were not as widely available. Their calculated values exhibited an average deviation of 2.31 % form the experimental data points considered in their study. A 1984 study by Sengers et al. [30] documents the available data on thermal conductivity of water since the promulgation of the first international formulation for transport properties of water substance in 1964. Their survey of experimental information summarized the thermal conductivity of water and steam from forty-three literature sources and only five sets were available for sub-ambient temperatures. About 40% of the datasets had been obtained using the transient hot-wire method. Another study by Nieto de Castro et al. [26] in 1986 examined the available thermal conductivity data for water with the purpose of establishing standard reference values along its saturation line. They noted that the thermal conductivity of fluids was one of the most difficult properties to measure and only after the technical advances during the 1970's the precision of these

measurements has significantly improved. Another set of standard reference data based on new experimental data was proposed by Ramires et al [31] nearly a decade later that led to a revised and more accurate correlation. Both the aforementioned studies [26, 31] fitted experimental data from multiple sources into a quadratic function of temperature. Ramires et al.[31] described the reduced thermal conductivity of water over the normal liquid range as $k^* = -1.48445 + 4.12292T^* - 1.63866T^{*2}$ for $274 \le T \le 370$ K.Where $T^* = T/298$ and $k^* = k(T)/k(298)$ with $k(298.15, 0.1 MPa) = 0.06065 \pm 0.0036$ $Wm^{-1}K^{-1}$ We will later compare the current experimental results to this correlation of Ramires et al. [31]. This study reports the thermal conductivity of water, ethylene glycol, glycerol, and propylene glycol.

2.2 Apparatus

The transient hot-wire apparatus used in the experiments consists of a stainless steel cylindrical cell 43 mm in diameter and 150 mm length. The cell is closed at the bottom and has a top lid with two electrical feedthroughs. Connected to these two electrical feeds are two copper conductors with tabs and screws to secure the platinum wire. The platinum wire is soldered to the copper conductor after securing it to the tabs to ensure a good contact. The platinum wire used in the measurement cell has a radius of 25 microns. Its purity is 99.99 % and it has is used in the hard drawn state. This wire acts as the line heat source. The length of the wire is estimated to be 95.33 mm. The length of the wire was obtained after the wire had been soldered in place by comparing its resistance to a known length of identical platinum wire. The resistance measurements were made in the 4-wire configuration using a Keithley 2440 sourcemeter. This sourcemeter was also used as a constant current source for heating the wire during the experiments. The terminal voltage at the feedthrough was measured using a 24 bit delta-sigma analog to digital converter with a nominal input voltage range of ± 10 V. The measuring cell

was maintained at a constant temperature by immersing it in a circulating bath such that only the top terminals were accessible for connections. A schematic of the test cell is shown in Fig. 2.1.



Figure 2.1: Schematic of the transient hot-wire cell.

The test cell is filled with technical grade distilled water such that it fully immerses the platinum wire. The cell is next immersed into the circulating bath and allowed to stabilize before starting measurements. The temperature of the test fluid is measured using a 1/16 inch K type thermocouple with exposed junction. This temperature was acquired by a national instruments module NI 9213, the accuracy of which is less than 0.02°C for high resolution mode, and less than 0.25°C for high speed mode. Also, the standard limits of error for the type K thermocouple in the range of 0 to 1250 °C is less than 2.2°C or 0.75%, and 2.2°C or 2 %.for -200 to 0°C.

The transient hot wire measurements consist of sending a 250 mA current pulse through the platinum wire and recording the terminal voltage using the 24 bit A/D converter. The sourcing

of the current and all other measurements are initiated using a LabView program developed specifically for this experiment. The voltage data is sampled at a frequency of 4167 Hz and 25,000 samples for a duration of 6 seconds are logged.

The temperature rise of the wire, when subjected to step current change, is determined by measuring the change in resistance with time. This change in resistance due to Joule heating is related to the temperature rise, and is given by [32]:

$$R(T) = R_{ref} \left[1 + \alpha_{ref} \left(T - T_{ref} \right) \right]$$
(2.5)

where the reference values correspond to the values at initial time (t = 0) prior to the passing the current. The value of R_{ref} is taken to be that of the first data point (t = 0) and the reference temperature corresponds to the temperature of the test liquid.



2.3 Determination of Temperature Coefficient of Resistance

Figure 2.2: Variation of platinum wire resistance with temperature.

Recognizing that we need to vary the reference temperature, it becomes necessary to first determine the temperature coefficient of resistance (α_{ref}) at each reference temperature. That is, the temperature coefficient of resistance needs to determine as a function of temperature. This is accomplished by a submerging the wire holder in a circulating bath and recording its resistance at varying bath temperatures. The resistance measurements are carried out in a 4-wire configuration with a small current (1 mA) that causes negligible heating in the wire. The variation in resistance of the wire as a function of temperature is shown in Fig. 2.2.

The local slope of the resistance versus temperature curve is next determined by selecting 75 consecutive points above and below the desired reference temperature. An example of the aforementioned procedure for obtaining the local slope in the neighborhood of T_{ref} =293.2 K is shown in Fig. 2.3.



Figure 2.3: Determination of local Slope.

A similar procedure is repeated at each reference temperature to obtain the dependence of α on temperature, the results of which are shown in Fig. 2.4.



Figure 2.4: Variation of the temperature coefficient of resistance for platinum wire with temperature.

Having determined the temperature coefficient of resistance, the temperature rise of the wire due to the passage of current can be calculated using equation (4). Note that we are now recording the temperature rise during the 6-second interval that corresponds to the passage of the current. Therefore, we assume that the very first deduced temperature data point corresponds (t = 0) corresponds to no heating of the wire, i.e. $\Delta T = T_{wire} - T_{liquid} = 0$. An example of the calculated temperature rise of the wire, for a 250 mA current, is shown in Fig. 2.5.



Figure 2.5: Temperature-time history for experimental runs at varying liquid temperatures. It can be seen from Fig. 2.5 that during the initial period of approximately 0 - 0.2 s the temperature rise exhibits a non-linear trend with respect to time on a logarithmic scale. This is followed by a linear rise in temperature with time on the log scale. A straight line is fit to the ΔT with ln(t) trace in the region of linear temperature rise, and its slope (S) is used in equation (1), which is $k = q/4\pi S$.

2.4 Data Reduction Procedure

The choice of the range of times (t_1, t_2) to use in the linear fit in order to obtain the slope S has generally been determined based on the experience of the experimenter. This is also noted in the work by Roder at the National Bureau of Standards [19] who states that operator judgement is involved in the selection of times t_1 and t_2 . We find that this is true for our experiments as well and propose a method to make the data reduction procedure more objective and free from experimenter bias. The theory suggests fitting a straight line in the region of interest ($Fo \gg 1$) assuming that the medium behaves as a solid. However, there is a practical problem with this in that the non-ideal effects are present both in the initial and terminal regions of the curve. They arise due to the finite heat capacity of the wire and the onset of convection at sufficiently large time. The existence of three distinct regions of increasing, apparently constant, and decreasing slopes due to these non-ideal effects is shown in Fig. 2.6 at a temperature of 320 K.



Figure 2.6: Illustration of the three distinct temperature rise regions observed in the experiments.

For the current work, as a first estimate, we choose the start time corresponding to the 1000th data point ($t_1 = 0.2398 s$) and the end time to be the last acquired data point ($t_2 = 5.9998 s$). This choice is based on a visual determination of the linear region observed in the experimental curve for ΔT versus time. An example of the range chosen for the fit is shown in Fig. 2.7a. The later onset of convection is however not apparent at lower temperatures of 274.5 K in the Fig. 2-7a. We find that a non-linear fit (3rd order polynomial) to region of interest helps us identify the regions of increasing and decreasing slope even in what visually appears to be straight-line segment for low temperature cases. For example, in Fig. 2-7a, the initial

curvature is obvious for -8 < ln(t) < -2, but the effect of the onset of convection is not readily obvious in the range -1.428 < ln(t) < 1.792. Performing a third order fit helps us identify the point where the slope begins to decrease at longer times. This point has been identified by a filled symbol in Fig. 2-7b. Note that the value of the peak slope for the polynomial fit is within 0.7 % of the linear slope for this specific case. The maximum slope for the polynomial fit allows for a consistent choice of a value for the slope, unlike the linear fit, where the value will depend on the start and end points chosen by the experimenter. We therefore use the peak value of the slope (S) obtained using a third order polynomial fit to the apparent linear region to estimate the thermal conductivity using $k = q/(4\pi S)$. Note that the region of interest (apparent linear) must still be carefully be chosen while performing the fit.



Figure 2.7: (a) Region of experiment used for fitting and extracting the slope (b) Comparison of slopes obtained using a linear and a 3rd order polynomial fit. The filled symbol represents the slope used in the experimental determination of thermal conductivity.

Chapter 3: Temperature Dependence of Thermal Conductivity for Water: Experiments and Computations

3.1 Thermal Conductivity Results for Water

Experimental thermal conductivity values for water for three different sets of experiments are shown in Figs. 3.1 (a)-(c) to illustrate the reputability and extent of scatter in the experiments. The three sets of tests were conducted with different samples of technical grade water and the temperature was slowly raised from approximately 273 K to 300 K over a 12-hour duration. The data show an increasing trend of thermal conductivity with an increase in temperature.



Figure 3.1: Three sets of experimental runs showing thermal conductivity variation with temperature for water.

It can be seen from Fig. 3.1 that the scatter in the experimental results tends to increase with rising temperatures. The three datasets are combined to obtain the overall results that show the temperature dependence of thermal conductivity of water. It is interesting to note that there is a local minimum in the thermal conductivity of water at around 276.89 K (3.74 °C). This slight dip appears in all the three data sets collected, as well as the consolidated data. We suspect that this is a related of the non-linear variation in density around this temperature range. Note that the maximum density for liquid water is close to 277.13 K (3.98 °C) [33].



Figure 3.2: Consolidated data set for thermal conductivity as a function of temperature for current experimental sets for water.

The residuals of the fit to the experimental data are shown in Fig. 3.3. The plot in Fig. 3.3 shows that the largest one-sided residuals are observed at around 4 degree Celsius. The confidence and prediction bands for the regression analysis are shown as well.



Figure 3.3: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data.

3.1.1 Comparison with Previously Reported Data

The current experimental results for the temperature dependence of thermal conductivity of water are compared to that reported in recent literature. The comparison also includes a dashed line showing the computed values obtained from the standard reference data from the study of Nieto de Castro et al.[26]. The work of Nieto de Castro et al. describes the reduced thermal conductivity of water as a second order polynomial over the temperature range of 274-370 K [26]. The plot in Fig. 3.4 shows that the current experiments consist of significantly larger number of data points over the temperature range of interest, exhibit the least scatter among all the experimental data sets, and are closest to the standard reference data. The current experimental observations are also nearly continuous over the entire temperature range. The

plot also illustrates the utility of the current work in that it extends the experimental data to lower temperature values close to the freezing point of water.



Figure 3.4: Comparison of current data with other recent studies for water.

3.2 Computed Results

3.2.1 Computational Specifications and Results

The heat transfer problem of an electrically heated platinum wire immersed in water is solved using the ANSYS Fluent CFD code. The problem is treated as axisymmetric because of the cylindrical symmetry. The governing equations for mass, momentum and energy conservation are given by [34]

$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho v_x)}{\partial x} + \frac{\partial (\rho v_r)}{\partial r} + \frac{\rho v_r}{r} = 0$$
(3.1)

$$\frac{\partial(\rho\vec{v})}{\partial t} + \nabla \cdot (\rho\vec{v}\vec{v}) = -\nabla p + \nabla \cdot (\bar{\tau}) + \rho\vec{g} + \vec{F}$$
(3.2)

$$\frac{\partial(\rho E)}{\partial t} + \nabla \cdot \left(\vec{v}(\rho E + p)\right) = \nabla \cdot (k\nabla T) + S_h \tag{3.3}$$

Where x and r represent the axial and radial coordinates. The stress tensor $\overline{\overline{\tau}}$ in the momentum equation is given by:[34]

$$\bar{\bar{\tau}} = \mu[(\nabla \vec{v}) + (\nabla \vec{v})^T - 2/3 \,\nabla \vec{v} l]$$
(3.4)

The quantities \vec{F} and S_h in the momentum and energy equation represent the external body force and volumetric heat source term, respectively. The current problem is treated as a natural convection problem, and the Boussinesq approximation is used for the density in the buoyancy term of the momentum equation.

$$\rho = \rho_o (1 - \beta (T - T_o)) \tag{3.5}$$

This approximation is valid as the temperature rise ΔT in the system is less than 3 Kelvin with $\Delta T \ll 1$. Here $\beta = -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial T}\right)$ is the volumetric thermal expansion coefficient, and ρ_o is the constant density of the flow at an operating temperature of T_o . The SIMPLE algorithm with a second order discretization scheme was used in the solution procedure.

Given the wide range of the length scales in the experimental apparatus ($a = 25 \times 10^{-6}m$, $L \approx 9.5 \ cm$) it is prohibitively expensive to compute the solution for the entire domain while resolving the flow and thermal features near the heated wire. A section of the wire corresponding to length $150 \times a$ is considered along the axial direction, with r = 0 being treated as the axis. The radial extent of the problem is limited to four times the section of the wire length i.e. $600 \times a$. The grid distribution used is shown in Fig. 3.5 (a). A non-uniform grid distribution with a bias factor of 5 was used in the radial direction for the fluid domain while a constant number of equally spaced divisions was set for the solid. The electrical heating of the
wire is represented by a volumetric heat source within the platinum wire. The strength of the volumetric heat source (Q) is obtained from the known experimental value of the power supplied (317.95 mW) and the wire dimensions ($a = 25 \ \mu m$; $L = 9.533 \ cm$). Gravity acts along the axial direction as indicated by the arrow in Fig. 3.5 (a). The Material properties used in the simulations are summarized in Table 3-1.

Property	Method	Value		
Water				
Density (kg/m3)	Boussinesq	9.987E + 02		
Cp (Specific Heat) (J/kg-K)	constant	4.186E + 03		
Thermal Conductivity (W/m-K)	constant	5.9374E - 01		
Viscosity (kg/m-s)	constant	1.0695E - 03		
Thermal Expansion Coefficient (1/K)	constant	1.7824E - 04		
Platinum				
Density (kg/m3)	constant	2.145E + 04		
Cp (Specific Heat) (J/kg-K)	constant	1.330E + 02		
Thermal Conductivity (W/m-K)	constant	7.160E + 01		

Table 3-1 Material Properties used in the simulations.

A pressure-inlet boundary condition was specified for the right and top boundaries with flow direction being normal to the boundary, while the left boundary was set to be of type pressureoutlet. The initial conditions correspond to a value of T = 290.53 K for both the fluid and solid domain, with a stationary fluid.

A solution corresponding to t = 6s is shown in Fig. 3.5 (b). This time corresponds closely to the last acquired data point in our experiments (Fo = 1363) when the effects of convection are quite apparent. The streamtraces in Fig. 3.5 (b) clearly show the buoyancy induced flow along with the contours of the temperature in the fluid as well as the solid at the end of t = 6 s. The maximum Rayleigh number based on the wire length, and the experimentally deduced $\Delta T = 2.87 K$ at t = 6s, is calculated to be $Ra_L = \frac{g\beta\Delta TL^3}{\nu\alpha} = 2.86 \times 10^7$, and as such, a laminar free convection flow is to be expected. The transition to turbulence generally occurs over the range of $10^8 < Ra < 10^{10}$ [34].



Figure 3.5: (a) A zoomed in view of the grid distribution. Note that the total simulated domain extends to 150 and 600 non-dimensional units in the axial (x/a) and radial(r/a) directions respectively. Here a is the radius of the wire and g the acceleration due to gravity. (b) Streamtraces and temperature contours at t=6.0 s.

For the given configuration, a characteristic reference velocity and can be defined as $u_0 = \sqrt{ag\beta(T_w - T_\infty)}$, where *a* is the wire radius. The wall temperature for evaluating this reference velocity is obtained at the axial midpoint of the wire. This velocity is can be used to non-dimensionalize the magnitude of the absolute velocities. Similarly, the temperature difference can be non-dimensionalized as $\lambda(T - T_\infty)/(Qa^2)$. Here *Q* is the strength of the volumetric heat source term in Watts per cubic meter.

3.2.2 Temperature Distribution and Local Heat Transfer Coefficients

The computed non-dimensional temperature and velocity magnitude contours corresponding to a time of 6-seconds is shown are Fig. 3.6 (a) and (b) respectively



Figure 3.6: Computed contours of non-dimensional (a) Temperature rise at t = 6s, and (b) the radial temperature distribution midway at x/a = 75 at various times.

The contours of the non-dimensional temperature rise field show that the temperature along the wire surface increases as one traverses the axial direction from the right boundary inlet against the direction of gravity. This is expected since the coldest fluid is in contact with the wire surface at the inlet (x/a = 150). Fig. 3.6 (b) also shows the time evolution of the radial temperature distribution at the axial location corresponding to the wire midpoint. There is a relatively large change in ΔT at a given radial location during the earlier times, in that the curves are widely spaced between t = 0.024 and t = 0.1 as compared to that between t = 4 and t = 5 seconds. An interesting feature of the radial temperature distribution is that at any given time the trend is mostly logarithmic (straight line on a log scale) with respect to the radial distance except at large distances. The slopes of the straight-line segments on log-plot shown in Fig. 3.6 (b) are also very similar. This seems to suggest that the heat transfer is dominated by conduction

and the fluid medium is apparently stationary. Furthermore, the radial extent diffusion of heat is limited to approximately 40 wire diameters.

Another important quantity of interest is the local heat transfer coefficient at the wire surface. The local heat transfer coefficient can be written as $h_{\chi} = \frac{q_{W}'(a,\chi)}{T_{W}(a,\chi) - T_{\infty}} = \frac{-k_{f} \frac{\partial T}{\partial r}|_{r=a}}{T_{W}(a,\chi) - T_{\infty}}$. Here k_f is the fluid thermal conductivity, $T_w(a, x)$ the local wire surface temperature, and T_∞ the farfield temperature. The local heat transfer coefficient scaled as $h_x 2a/k$ is shown in Fig. 3.7 (a) as a function of the axial co-ordinate at the instant corresponding to 6-seconds along with the steady-state value with dotted lines for comparison. Note that a steady solution exists because of the nature of the boundary conditions imposed. The right and top boundaries have been specified as pressure inlet while the left boundary is a pressure outlet. This would not hold true for an enclosed container, but is appropriate for an infinite fluid medium. In order to examine the variation of the local heat transfer coefficient with the distance from the 'entrance' on the right boundary of the simulation we define x' = 150a - x, where 150a is the total axial extent of the simulation domain. This x' is consistent with the distance commonly used in the literature. A plot of the variation of the scaled heat transfer coefficient under steady-state conditions is shown with x'/a in Fig. 3.7 (b). The local heat transfer coefficient for the steady state solution exhibits a power law variation with respect to the distance from the entrance. Note that the computed data shown in Fig. 3.7 (b) begins and ends at x'/a = 1 and 149 respectively, i.e. one wire diameter has been excluded from either end.



Figure 3.7: (a) Scaled local heat transfer coefficient (a) as a function of distance (b) Fit (solid line) to the steady state distribution as a function of distance, x', from the entrance.

The variation of the scaled local heat transfer coefficient with time at a fixed location axially midway is shown in Fig. 3.8. The plot in Fig. 3.8 shows that the magnitude of the local heat transfer coefficient midway along the wire (x/a = 75).



Figure 3.8: Scaled local heat transfer coefficient at x/a = 75 as a function of time. Its magnitude decays monotonically with time and approaches a value of 0.3996 for the steady solution. The rate of change of the local heat transfer coefficient with time at the centerline is extremely rapid for Fo < 1, gradually tapers off in the range of 1 < Fo < 100, and is negligible for Fo > 1000. Note that the appearance of the rapid change of slope has been somewhat suppressed due to the stretching out of the time axis on account of being plotted on a logarithmic axis. The smallest and largest values of times in the plot correspond to 0.24×10^{-3} and 6 seconds, respectively.

3.2.3 Velocity Field

The velocity magnitude field for t = 6 seconds is shown in Fig. 3.9 (a). The contours correspond to values of $\frac{|V|}{\sqrt{ag\beta(T_w - T_{\infty})}}$. As noted previously, the wall temperature was evaluated at the axial midpoint on the surface of the wire. The peak velocity magnitude occurs radially at a distance of approximately 5 wire diameters from the surface and is shifted slightly toward the outlet side with respect to the center of the wire.

The peak velocity magnitude in Fig. 3.9 (a) is equal to 11.02 microns/second. This implies that a fluid parcel near the wire can translate roughly 0.441 wire diameters every second because of the buoyancy-induced force resulting from the density gradient. The axial velocity profile as a function of the radial coordinate is plotted for various times in Fig. 3.9 (b). The profiles were obtained at a fixed axial location in the middle of the domain. The axial velocity profiles exhibit a peak and then gradually reduce to zero within about 40 wire diameters.



Figure 3.9: Non-dimensional radial temperature and axial velocity profiles at x = L/2*.*

3.3 Comparison of Experimental and Computed Results

While there have been several experimental efforts using the transient hot-wire method that are far less studies that computationally examine the temperature and flow-fields for this configuration. The computational study shown in the previous sections enables us to compare the experimentally deduced temperature profiles to those obtained from simulations as shown in Fig. 3.10 (a). The agreement between the experimental and computed temperature is within approximately 0.5 K within the entire time span. The difference in temperature rise (ΔT) increases for up to approximately 0.01s and then is nearly constant for the subsequent time duration. Furthermore, the shape of the experimental temperature rise on a log scale during the early part of the transient t < 0.01 s. The slopes of the experimental and computed results are also very similar during later times.



Figure 3.10: Comparative temperature-time history for (a) Experiments versus CFD results, and (b) Exact solution to heat diffusion equation for a line source in an infinite media versus CFD results.

The temperature rise from the CFD simulations is also compared to the analytical solution for the heated line source in an infinite solid medium, as shown in Fig. 3.10 (b). It can be seen that there is a very good agreement between the two, thereby implying that there is only a limited influence of the convective effects on the temperature rise for the duration of heating under consideration.

3.4 Summary

The transient hot-wire technique was used to obtain the temperature dependence of the thermal conductivity of liquid water over a temperature range of 273 to 301 K. The experimentally determined values show close agreement to the standard reference values. We find that in order to obtain the correct temperature dependence of the thermal conductivity it is essential to determine the variation of temperature coefficient of resistance of the platinum wire and not assume it to be a constant. A method to minimize user bias in determining the slope of the temperature-rise versus logarithm of time curve used to determine the thermal conductivity was presented. 2-D Axisymmetric CFD simulations were conducted to understand the nature of the

temperature of velocity field in the vicinity of the wire. A comparison of the experimental and computed results showed a good agreement for the temperature rise of the wire with respect to time at a constant heating rate.

Chapter 4: Temperature Dependence of Thermal Conductivity for Ethylene Glycol, Glycerol, and Propylene Glycol

4.1. Ethylene Glycol

Ethylene glycol is an organic compound produced on an industrial scale. The molecular formula of ethylene glycol is $C_2H_6O_2$, and the melting and boiling point are 260.5 *K* and 270 *K* respectively [35, 36] The compound is a common heat transfer fluid and used as an antifreeze and deicing agent. It is also used as a solvent and in hydraulic brake fluids. It is therefore important to study the thermal transport properties of this liquid.

DiGuillo and Teja [35] were two early experimentalists who measured the thermal conductivity of the first-six members of the poly ethylene glycols and their binary blends. Their study also included measurements for ethylene glycol. The temperature range of their experiment was from 298.6 K - 471.3 K, and the reported thermal conductivity for ethylene glycol was $0.2541 W/(m \cdot K)$ to $0.2444 W/(m \cdot K)$ at the high and low temperature limits. They used the transient hot-wire method to carry out their measurements. Their results exhibit a slight increase in the thermal conductivity of ethylene glycol with temperature for temperatures between 290 – 415 K followed by a reduction for temperatures up to 480 K.

Another study by Azarfar et al. [37] also used the transient hot-wire to study the thermal conductivity of ethylene glycol. They used a copper micro-wire instead of the conventional platinum wire in their apparatus. Measurements for ethylene glycol were carried out in the temperature range of 283 - 313K. Their reported thermal conductivity values for the thermal conductivity of ethylene glycol were 0.2433 and 0.2537 $W/m \cdot K$ for temperatures of 283 and 313 *K*, respectively. They reported that the average uncertainties were in the range of ± 0.9 percent.

Lin et al. [38] used molecular dynamics simulations to investigate the thermal conductivity of ethylene glycol. This study computationally identified the major contributions to thermal conductivity from rotational energy transfer, intramolecular interactions, and hydrogen bonds. They conclude that intramolecular hydrogen bonds can have a major influence on the variation of thermal conductivity with temperature. A temperature range of 298 - 470 K was investigated in their work. Various molecular models were used to obtain thermal conductivity as a function of temperature. The thermal conductivity was relatively high at 298 *K* compared to Digullo's and Teja's [35] experiment, but the thermal conductivity showed a good match at 470 *K*.

Beck et al. [39] enhanced the thermal conductivity of ethylene glycol by adding aluminum oxide nanoparticles to ethylene glycol. They measured the thermal conductivity by using the hot-wire transient method, and the temperature range was 290 to 298 411 K. The volume fraction of Al_2O_3 was 1% to 4%, and the thermal conductivity exhibited and increase to 0.285 $W/(m \cdot K)$ at 409 K in 3.71% of the volume fraction of Al_2O_3 .

For the current experiments presented in this thesis, ethylene glycol was purchased from Fisher Scientific, and the purity of ethylene glycol is greater than 99.96%. In this current work, the thermal conductivity of ethylene glycol was measured over a temperature range of 260 - 340 K using the transient hot-wire method. The temperature range was traversed thrice to produce three data sets which were the aggregated to obtain the final results for the temperature dependence of the thermal conductivity.

4.1.1 Thermal Conductivity Results for Ethylene Glycol

The results for three different experimental runs for ethylene glycol are shown in Figs. 4(a)-(c). It can be seen that the thermal conductivity of the three sets shows an increasing trend with

temperature. For the current ethylene glycol experiments, more than 300 pieces of individual data are collected for each run. It was found that for each run, convection starts at around 340 to 350 K. Therefore, the data collection was stopped in that temperature range. There is an excellent run-to-run reproducibility among the three datasets shown in Fig. 4.1.



Figure 4.1: Three sets of experimental runs showing thermal conductivity variation with temperature for ethylene glycol.

The three datasets are combined to obtain a consolidated plot for the thermal conductivity as a function of temperature and is shown in Fig. 4.2. A linear fit to the data is also shown as a dashed line. It can be seen that overall the thermal conductivity shows an increasing trend with temperature for the given test range. The slope of the straight line for ethylene glycol $(b = 1.8548 \times 10^{-4})$ though is significantly lower compared to that of water $(b = 18.482 \times 10^{-3})$. The thermal conductivity variation for ethylene glycol is approximately 6.16 % over the entire temperature range.



Figure 4.2: Consolidated data set for thermal conductivity as a function of temperature for current experiments for ethylene glycol.

The experimental results a 6.2 percent change in thermal conductivity when the temperature varies over the range of 260 - 340 K.

For clarity, Fig. 4.3 separately plots the residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for ethylene glycol. There are very few data points that are outliers when assessed by the criterion of inclusion within the prediction band.



Figure 4.3: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for ethylene glycol.

In order to validate the current results we compare it with data reported in some recently published work. Figure 4.4 shows the comparison between the current work, the study Azarfar et al. [37], Digullo et al. [35], and Pastoriza-Gallego et al. [40]. The current results are in good comparison with all three previously reported work. Interestingly, the current experimental results fall between the results of DiGuilio et al. [35], and Pastoriza-Gallego et al. [40] over the entire temperature range, with the results from Pastoriza-Gallego et al. [40] being consistently lower than the current study. The results for the other two studies are slightly on the higher side as compared to the current data. Furthermore, though the results of Pastoriza-Gallego et al. [40] exhibit as slope similar to the current data for temperatures below 323 *K*. Another noteworthy point is that all the datasets, including that of the current work, show an increase in thermal conductivity of ethylene glycol with increasing temperature. Note that that the comparison in Fig. 4.4 has been restricted to the first of the three runs in the current study. This is to enable

clarity in the Fig. 4.4, as the combination of all three sets would suppress the limited number of data points from the other studies.



Figure 4.4: Comparison of current data with other recent studies for ethylene glycol.

4.2 Glycerol

The word glycerol derives from the Greek term for sweet "glykys", because of its sweet taste. Glycerol was discovered in 1779 by the Swedish chemist Carl W Scheele. The molecular formula of glycerol is $C_3H_5(OH)_3$. It has a melting point close to room temperature (291.35 *K*), and the boiling point is 563.15 *K* [41, 42]. It is also commonly referred to as glycerin or glycerine. Glycerol-water solution is used to prevent freezing at low temperatures, and it was used as antifreeze before the discovery of ethylene glycol. At low temperatures, glycerol has a tendency to supercool rather than crystallize [43]. Given its interesting heat transfer characteristics it was decided to study the response of its thermal conductivity to temperature variation at temperatures near and above room temperatures.

Among the studies related to the measurement of transport properties of glycerol are those of Sun et al. [44] who used a laser based thermal pulse technique to make simultaneous measurements of thermal conductivity and diffusivity. There experimental efforts yielded a value of $0.29 W/(m \cdot K)$ corresponding to a temperature value of 296.45 K. The principle of their measurements was based on photothermal deflection of a laser probe beam that occurs due to a temperature gradient resulting from a prior square heating pulse applied to a thin heating wire. This work provided values for thermal conductivity at a single temperature. Determination of the thermophysical/transport properties in their work was done by fitting computed results for solutions to the heat conduction problem to the measured time-dependent beam deflection.

Another experimental study that reported the thermal conductivity of glycerol is that by LeBrun and Markides [45] who used deionized water and glycerol as validation targest for a Galinstan-filled capillary probe for thermal conductivity measurement of molten salts. Their apparatus was a modified transient-hot wire and was designed to function in harsh environments and work with ionic liquids under high temperature conditions. The sensing element in their device comprised of a U-shaped quartz-capillary filled with liquid Galinstan. The temperature range of their measurements for the thermal conductivity of glycerol was between 301.7 - 431.2 K, and the thermal conductivity at those two temperatures were $0.2903 W/(m \cdot K)$ and $0.3028 W/(m \cdot K)$, respectively. Their reported estimate of error for the measurements was less than 2%.

There have also been attempts to enhance the thermal conductivity of heat transfer liquids by dispersing nanoparticles in them in order to create so-called nanofluid. A study by Sharifpur et al. [46] looked at enhancing the thermal conductivity of glycerol by mixing glycerol with nanoparticles of aluminum oxide. They measured the thermal conductivity with nanoparticles of three different sizes corresponding to 31 nm, 55 nm and 134 nm in diameter. They pointed out that the thermal conductivity of the glycerol-based nanofluids with aluminum oxide Al_2O_3 increased by 19.5% at a 4% volume fraction for the 31 nm particles. The device used in their measurements was from a commercial vendor (KD2 Pro, Decagon Devices, Pullman, WA) and the temperature range reported in their study was from 293.15 – 318.15 *K*. The underlying principle of KD2 Pro was the transient hot-wire method. The reported uncertainty in measurements of the thermal conductivity was between 5.1% – 8.5%.

Another effort to increase the thermal conductivity of glycerol was reported by S. Akilu et al. [47] who dispersed silicon carbide nanoparticles in a mixture of glycerol and ethylene glycol The ratio of the glycerol to ethylene glycol was in the proportion of 60: 40 by weight. The temperature range of their experiment was between 288.15 - 348.15 K, and the size of the *SiC* nanoparticles of ranged between 45 - 65 nm. The maximum concentration of nanoparticles in the mixture was 1% of weight. They noted that the thermal conductivity of the mixture with the nanoparticles of (*SiC*) increased 24.5% at 304 K compared to the base liquid mixture.

The current experiments examine the thermal conductivity of neat glycerol in the temperature range of 290 - 350 K. Glycerol for this study was obtained from Macron Fine ChemicalsTM, and had a purity of 99.99%. The temperature dependence was obtained by concatenating three independent sets of data that were collected during different days over a one-month period. Note that the test for glycerol begin very close to the room temperature since it has a melting point of 291.35 *K*.

4.2.1 Thermal Conductivity Results for Glycerol

The plots in Fig. 4.5 (a-c) show the results of the three sets of experiments conducted to obtain the thermal conductivity of glycerol. There is a relatively larger scatter in the data compared to ethylene glycol over the same temperature range. One must not over-interpret the result since the variation at a given temperature is still less than $0.003 W/m \cdot K$. The plots show a very good run-to-run reproducibility with very similar thermal conductivity values both the lower and higher temperature ends.



Figure 4.5: Three sets of experimental runs showing thermal conductivity variation with temperature for glycerol.

The data points obtained from the three distinct runs are combined into a single plot. A linear regression is performed on the consolidated data set. The slope of the line fit to thermal conductivity versus the temperature for glycerol ($b = 1.1321 \times 10^{-4}$) is of the same order of magnitude, but slightly lesser that of ethylene glycol ($b = 1.8548 \times 10^{-4}$). The thermal conductivity variation for glycerol is approximately 3.34 % over the entire temperature range while a 6.16 % was observed for ethylene glycol for the same temperature range.



Figure 4.6: Consolidated data set for thermal conductivity as a function of temperature for current experimental sets for glycerol.

A most remarkable point is that the current apparatus can distinguish very minute differences in thermal conductivity with relative ease. This can be clearly seen if one compares the values for thermal conductivities for glycerol with ethylene glycol as predicted by the straight line fit for the two liquids at, say 300 K. At that temperature the difference in the predicted thermal conductivities differ by only 0.02752 units.

For clarity, Fig. 4.7 separately shows the residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for glycerol. Most of the outliers are seen to occur in the high temperature range with data points located outside the upper line for the prediction band.



Figure 4.7: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for glycerol.

The current experimental data is next compared to some previously reported results in the literature. The comparative plot in Fig. 4.8 shows that the study of Sharifpur [46] shows a good agreement in the low temperature region but deviates towards higher values for higher temperatures. Another set of data from the study of Gelder [48] is higher by approximately 4% over the entire temperature range. Finally, the results reported by LeBrun and Markides [45] are the highest among all the reported results. These differences, except for the results by LeBrun and Markides [45] are less than 10%. Therefore, we can say that the current experiments are in fair agreement with most of the previous studies. Note however, that all the three datasets predict an increasing thermal conductivity for glycerol with increasing temperature, similar to the current results.



Figure 4.8: Comparison of current data with other recent studies for glycerol.

4.3 Propylene Glycol

Propylene glycol is a synthetic liquid which is hygroscopic in nature. It is also known as such as 1,2-dihydroxypropane, and 1,2-propanediol. The molecular formula of propylene glycol is $C_3H_8O_2$. Propylene glycol is used in producing polyester compounds, de-icing solutions, antifreeze and in heat transfer the same way ethylene glycol is used. For safety reasons, it is preferable to use propylene glycol instead of ethylene glycol because of the high toxicity of ethylene glycol [49]. The melting point of propylene glycol is 213.15 *K* and the boiling point is 460.15 *K* [50] Sun and Teja [51] studied the thermal conductivities, densities, and viscosities of propylene glycol, dipropylene glycol, and tripropylene glycol. The technique used in their experiment was the transient hot-wire method. The temperature range for their experiment was 290 K - 460 K. The reported uncertainty of their results was $\pm 2\%$.

Deng et al. [52] examined the thermal conductivities of 1, 2-Ethanediol and 1, 2-Propanediol also using the transient hot-wire method, and the temperature range for the experiments was 253.15 - 373.15 K, at atmospheric pressure condition. The uncertainty in the results for thermal conductivity was better than $\pm 2\%$. Interestingly, their results showed that the thermal conductivity decreased with increasing temperature.

Another study by Palabiyik et al [53] was related to the enhancement of the thermal conductivity of propylene glycol by the addition of Al_2O_3 and TiO_2 nanoparticles. The temperature range for this study was 293.15 – 353.15 K. The size of nanoparticles of Al_2O_3 and TiO_2 nanoparticles were 13 nm and 21 nm, respectively. They compared their result for thermal conductivity of pure propylene glycol with Sun's and Teja's work [51], and found that the difference between their and Sun's and Teja's results was less than ± 5%. They reported that the thermal conductivities of propylene glycol based Al_2O_3 and TiO_2 nanofluids showed an increase of 11%, and 9%, respectively.

Suganthi et al. [54] obtained the thermal conductivity of propylene glycol based ZnO nanofluids as a function of temperature and nanoparticles concentration. The temperature range for their study was 283.15 – 333.15 *K*. The volume fraction of nanoparticles was less than 2%. The size of nanoparticles was between 35 and 40 nm. They report that the highest enhancement of the thermal conductivity occurred at the lowest temperature.

We examine the thermal conductivity of pure propylene glycol in this study in the temperature range of 235 - 360 K. Propylene glycol was purchased from Tokyo Chemical Industry Co., Ltd, and its purity was greater than 99.0 %.

4.3.1 Thermal Conductivity of Propylene Glycol

In a procedure similar to the other three liquids, we obtain three independent data sets over a wide range of temperature, which are shown in Fig. 4.9. The three are very similar to each other. The only difference for these three sets is that the readings were collected while the temperature range was traversed from the high to the low temperature region. The data for the three other fluids presented previously were collected with the temperature range being traversed from the low to the high region. Regardless of the direction of traversal, the propylene glycol thermal conductivity still shows an increase with increasing temperature. The experiments had to be limited to a value of 360K because of the onset of convection for the current operating conditions. Typically, about 180 thermal conductivity data points were collected over the entire temperature range for each of the three runs.





Figure 4.9: Three sets of experimental runs showing thermal conductivity variation with temperature for propylene glycol.

The three sets of data shown in Fig. 4.9 are combined to produce a consolidated data set. A linear regression is done on the combined data set to assess the influence of temperature by examining the slope of the fitted thermal conductivity values with respect to temperature.



Figure 4.10: Consolidated data set for thermal conductivity as a function of temperature for current experimental sets for propylene glycol.



Figure 4.11: Residuals, confidence, and prediction bands (95%) for the linear fit to the current experimental data for propylene glycol.

The slope of the fit for propylene glycol lies midway between those of ethylene glycol and glycerol. It shows a 4.89 % variation in thermal conductivity between the temperatures of 260 - 340 K as determined from the coefficients of the linear fit. The linear fit for the entire propylene glycol data, and the residuals, confidence, and prediction bands (95%) are shown in Fig. 4.10 and 4.11, respectively.

A comparison of the current experimental data set with previous studies in the literature for propylene glycol is shown in Fig. 4.12. The experiment results of Cabaleiro et al.[55] are about 8% higher than the current study, but show an increase with temperature similar to this work. On the other hand, the experimental of Sun et al.[51], and the experiment of Suganthi et al.[54] show that the thermal conductivity decreases when the temperature increases. Note however that both their results are in good agreement with the current study in the high temperature region. This inconsistency related to the trend of variation of a materials thermal conductivity,

especially for liquids can be commonly found in the literature. This illustrates the importance of developing further systematic experimental and computational approaches to examine transport property trends with temperatures especially for fluids commonly using in engineering heat transfer applications.



Figure 4.12: Comparison of current data with other recent studies for propylene glycol.

4.4 Comparative Thermal Conductivities for the Four Liquids

A summary plot of the thermal conductivity of all of the four heat transfer liquids is shown in Fig. 4.13. The thermal conductivities for the four liquids can be ranked in the following order:

Water > Glycerol > Ethylene Glycol > Propylene Glycol

It can be seen that the temperature range of the results for water is significantly smaller than the other liquids. This is because it is extremely difficult to measure thermal conductivities for low viscosity liquids in general. Water also shows the largest variation in thermal conductivity even over this limited range. Another important conclusion that can be drawn is that the apparatus

and techniques developed in this work are capable of consistently resolving small differences between liquids such as glycerol and ethylene glycol. This minute difference is of the order of the experimental uncertainty and scatter that are found in data reported in the literature (cf. comparison plots with literature data for the fluids)



Figure 4.13: Comparative thermal conductivities of all four liquids obtained in this work. Having described the capabilities of the apparatus developed and the results obtained, it is also important to provide a measure of the variability in the current data set. We chose a small temperature interval 297 < T < 299 K, centered on 298 K to provide the extent of scatter. The bounds of the boxes indicate the limits of 25th and 75th percentile with the line as the median. The whiskers from the top and bottom of each extend to values that are 1.5 times the interquartile range, and the outliers are shown as empty circles.



Figure 4.14: Plot showing the variability in the thermal conductivity values for the four liquids in the neighborhood of $298 \pm 1K$.

Chapter 5: Conclusions and Future Work

5.1 Summary of this Study

A transient hot-wire apparatus was used to obtain the thermal conductivities of four common heat transfer liquids. A transient hot-wire apparatus was designed and constructed to enable these measurements. The cell used a 25 micron radius platinum wire that was 95.33 mm long. In addition to the cell, the setup required the integration of heating and cooling baths, highspeed data acquisition system, and the processing software. It was found that in order to obtain reliable estimates of the variation of the thermal conductivities with temperature it is important to account for the variation of temperature coefficient of resistance of the platinum heating/temperature sensing wire.

This study used a constant current of 250 mA for all the fluids studied, and the test time was limited to six seconds. A data reduction procedure for the unambiguous identification of the slope and the portion of the ΔT vs ln(t) was outlined. It was found that for the lowest viscosity fluid, water, effects of natural convection led to a non-linear relationship between the temperature rise and the logarithm of time after a time of approximately six seconds for the given current and wire length. Based on the experience for measuring the four different liquids it can be concluded that the measurement of thermal conductivity of low viscosity liquids is an extremely challenging task, especially as the temperature is raised. An earlier onset of convection occurs as the temperature is raised.

Liquid water was used for validating the results from this test cell, and the results showed a very agreement with the recommended values over the entire temperature range. Numerical studies were carried out to understand the effect of the transient heating of the wire on the temperature and velocity field in the surrounding medium. It was found that the computationally observed temperature rise of the wire agreed quite well with the theoretical results for a line source in an infinite medium, as well as the experimental data. This implies that there is only a limited influence of the convective effects on the temperature rise for the duration of heating under consideration. The effect of wire heating on the velocity field over a heating duration of six seconds (similar to the experiments) was found to be negligible beyond 40 wire diameters for the computations.

The temperature dependence of the thermal conductivities of water, ethylene glycol, glycerol, and propylene glycol were measured. All four liquids exhibited an increase in their respective thermal conductivities with increasing temperature. As previously noted, it was difficult to measure thermal conductivity of water above 300 *K* on account of its low viscosity. However, good experimental data for water was obtained from near its freezing point up to room temperature. The high temperature limit was approximately 340*K* for this study.

5.2 Future Work

There are several possibilities of improvement and additional work that could enhance this research. It may be worthwhile to obtain improved measurements of the local slope for the temperature coefficient of resistance as a function of temperature. Furthermore, these experiments could be extended to measure the thermal diffusivity of the liquids by suitable modifications to the apparatus. There is need for further research on explaining the differences observed in the temperature rise profiles between the experimental and computed results. Two specific suggestions for future work include (a) using the current apparatus to examine the effect of nanoparticle addition to the base fluid, and (b) characterization of thermal conductivity of phase change materials for both solid and liquid phases.

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	Ther	mal Con	ductivity	y V	alues fo	r Water	' in	W/m.K	(R	Run	-1)	
Т (К)	k	Т (К)	k		Т (К)	k		T (K)	k		T (K)	k
275.14	0.5676	285.37	0.5809		290.79	0.5908		296.08	0.6013		299.39	0.6034
275.29	0.5667	285.45	0.5832		290.87	0.5982		296.26	0.5979		299.42	0.6046
275.69	0.5645	285.53	0.5843		290.96	0.5921		296.30	0.6014		299.71	0.6117
275.91	0.5635	285.65	0.5819		291.13	0.5933		296.36	0.6069		299.84	0.6097
276.16	0.5641	285.82	0.5802		291.28	0.5890		296.48	0.5981		299.96	0.6054
276.97	0.5625	285.89	0.5822		291.33	0.5901		296.54	0.6002		299.96	0.6067
277.86	0.5625	286.00	0.5856		291.50	0.5931		296.58	0.5998		300.07	0.6099
278.02	0.5647	286.10	0.5847		291.61	0.5911		296.69	0.5977		300.09	0.6078
278.16	0.5624	286.26	0.5900		291.67	0.5914		296.79	0.6001		300.14	0.6155
278.35	0.5636	286.49	0.5830		292.03	0.5910		296.82	0.6027		300.19	0.6062
278.52	0.5667	286.63	0.5877		292.27	0.5959		296.98	0.6009		300.27	0.6143
278.65	0.5647	286.91	0.5849		292.31	0.5947		297.03	0.6059		300.28	0.6092
278.95	0.5683	287.00	0.5842		293.20	0.5922		297.15	0.6053		300.33	0.6112
279.41	0.5669	287.05	0.5824		293.32	0.5976		297.23	0.6102		300.40	0.6068
280.11	0.5701	287.26	0.5843		293.47	0.5987		297.33	0.6009		300.47	0.6117
280.42	0.5740	287.36	0.5825		293.50	0.5954		297.40	0.6079		300.49	0.6110
280.53	0.5769	287.44	0.5896		293.67	0.5954		297.66	0.6035		300.52	0.6069
280.65	0.5707	287.62	0.5847		293.72	0.6006		297.73	0.5994		300.57	0.6115
280.77	0.5737	287.71	0.5854		293.76	0.5944		297.85	0.6036		300.64	0.6049
281.02	0.5716	287.75	0.5861		293.82	0.5958		297.86	0.6013		300.69	0.6056
281.13	0.5721	287.95	0.5865		293.88	0.5956		297.92	0.6014		300.72	0.6082
281.25	0.5735	287.98	0.5892		294.05	0.5970		298.09	0.6016		300.80	0.6082
281.69	0.5745	288.07	0.5914		294.22	0.6019		298.15	0.6053		300.84	0.6108
281.94	0.5768	288.20	0.5868		294.22	0.6019		298.18	0.6027		300.86	0.6138
282.05	0.5738	288.28	0.5885		294.46	0.5935		298.33	0.6059		300.90	0.6110
282.65	0.5784	288.36	0.5898		294.52	0.5947		298.38	0.6045		300.93	0.6125
282.76	0.5751	288.45	0.5862		294.58	0.5966		298.44	0.6086		301.00	0.6097
283.17	0.5772	288.58	0.5863		294.84	0.5960		298.45	0.6010		301.06	0.6074
283.34	0.5804	288.99	0.5882		294.87	0.5980		298.54	0.6038			
283.46	0.5814	289.26	0.5880		294.94	0.5963		298.62	0.6025			
283.54	0.5779	289.42	0.5880		295.22	0.6000		298.64	0.6075			
283.62	0.5805	289.60	0.5899		295.26	0.5995		298.69	0.6037			
283.92	0.5808	289.74	0.5882		295.63	0.5984		298.78	0.6037			
284.25	0.5779	289.94	0.5914		295.73	0.6035		298.83	0.6048			
284.32	0.5817	290.09	0.5865		295.80	0.5994		298.94	0.6074			
284.89	0.5839	290.18	0.5892		295.82	0.5965		298.96	0.6062			
285.08	0.5801	290.37	0.5937		295.89	0.5978		299.16	0.6115			
285.16	0.5809	290.53	0.5919		295.93	0.5984		299.28	0.6074			
285.27	0.5854	290.73	0.5944		296.03	0.5968		299.34	0.6064			

Appendix A: Tabulated Thermal Conductivity Results for Water
	Th	ern	nal Cono	ductivit	y V	alues fo	r Water	'n	W/m.K	(F	Run	-2)	
Т(К)	k		Т (К)	k		Т (К)	k		Т (К)	k		Т (К)	k
271.72	0.5577		275.19	0.5623		279.11	0.5651		282.60	0.5754		285.76	0.5824
271.66	0.5576		275.26	0.5595		279.20	0.5637		282.65	0.5760		285.80	0.5808
271.61	0.5558		275.33	0.5660		279.28	0.5669		282.72	0.5745		285.86	0.5860
271.60	0.5592		275.45	0.5632		279.36	0.5670		282.82	0.5779		285.92	0.5800
271.60	0.5584		275.54	0.5641		279.42	0.5642		282.87	0.5708		286.03	0.5794
271.62	0.5568		275.64	0.5625		279.50	0.5650		282.94	0.5738		286.06	0.5814
271.76	0.5568		275.73	0.5600		279.58	0.5704		283.00	0.5742		286.13	0.5814
271.93	0.5633		275.83	0.5599		279.67	0.5653		283.15	0.5772		286.21	0.5796
271.98	0.5634		275.90	0.5631		279.75	0.5719		283.24	0.5765		286.26	0.5804
272.07	0.5609		275.98	0.5592		279.82	0.5677		283.28	0.5758		286.42	0.5880
272.16	0.5597		276.06	0.5581		279.88	0.5686		283.33	0.5750		286.43	0.5823
272.33	0.5602		276.14	0.5593		279.97	0.5666		283.39	0.5777		286.49	0.5808
272.43	0.5608		276.22	0.5612		280.04	0.5662		283.46	0.5760		286.56	0.5814
272.51	0.5554		276.29	0.5598		280.15	0.5702		283.49	0.5784		286.60	0.5842
272.71	0.5579		276.36	0.5597		280.21	0.5711		283.58	0.5746		286.64	0.5809
272.81	0.5541		276.42	0.5621		280.37	0.5656		283.71	0.5730		286.69	0.5843
272.91	0.5600		276.50	0.5622		280.43	0.5726		283.80	0.5759		286.77	0.5865
273.02	0.5577		276.57	0.5635		280.48	0.5688		283.99	0.5755		286.85	0.5833
273.11	0.5609		276.63	0.5652		280.55	0.5723		284.06	0.5788		286.96	0.5786
273.20	0.5605		276.69	0.5623		280.71	0.5729		284.12	0.5815		287.01	0.5833
273.28	0.5623		276.78	0.5604		280.88	0.5729		284.17	0.5798		287.08	0.5854
273.35	0.5583		276.85	0.5599		280.91	0.5741		284.21	0.5794		287.18	0.5844
273.42	0.5574		276.91	0.5609		281.06	0.5687		284.26	0.5822		287.29	0.5820
273.54	0.5602		277.12	0.5614		281.13	0.5714		284.33	0.5786		287.34	0.5814
273.66	0.5624		277.23	0.5605		281.29	0.5694		284.59	0.5802		287.37	0.5818
273.74	0.5565		277.62	0.5610		281.34	0.5694		284.67	0.5784		287.42	0.5837
273.80	0.5628		277.68	0.5617		281.42	0.5719		284.75	0.5783		287.47	0.5869
273.92	0.5608		277.78	0.5594		281.50	0.5753		284.89	0.5813		287.51	0.5895
274.02	0.5585		277.89	0.5644		281.55	0.5741		284.95	0.5807		287.57	0.5888
274.11	0.5621		277.98	0.5623		281.71	0.5761		285.02	0.5802		287.63	0.5831
274.31	0.5630		278.16	0.5654		281.75	0.5720		285.08	0.5802		287.73	0.5840
274.38	0.5609		278.24	0.5642		281.84	0.5720		285.18	0.5791		287.82	0.5813
274.52	0.5606		278.33	0.5660		281.91	0.5732		285.26	0.5811		288.02	0.5879
274.58	0.5619		278.50	0.5648		282.14	0.5771		285.38	0.5823		288.06	0.5826
274.65	0.5647		278.59	0.5635		282.19	0.5707		285.47	0.5841		288.11	0.5862
274.74	0.5609		278.66	0.5662		282.34	0.5698		285.55	0.5815		288.31	0.5865
274.83	0.5627		278.80	0.5658		282.40	0.5725		285.60	0.5797		288.37	0.5823
274.91	0.5605		278.89	0.5659		282.45	0.5735		285.67	0.5816		288.43	0.5868
275.11	0.5627		279.02	0.5666		282.55	0.5761		285.69	0.5819		288.49	0.5876

	Th	ern	nal Conc	ductivit	y V	alues fo	r Water	' in	W/m.K	(R	lun	-2)	
T (K)	k		Т (К)	k		Т (К)	k		Т (К)	k		T (K)	k
288.52	0.5837		290.93	0.5930		292.96	0.5911		294.65	0.5959		296.76	0.6009
288.60	0.5911		291.06	0.5911		292.96	0.5961		294.73	0.5941		296.85	0.6029
288.67	0.5874		291.09	0.5907		293.00	0.5944		294.76	0.5930		297.13	0.6012
288.74	0.5877		291.14	0.5926		293.13	0.5974		294.79	0.5982		297.39	0.6036
288.78	0.5869		291.18	0.5906		293.16	0.5947		294.86	0.5935		297.50	0.6027
288.81	0.5880		291.24	0.5907		293.25	0.5966		294.88	0.5918		297.80	0.6028
288.85	0.5860		291.27	0.5881		293.27	0.5876		294.92	0.5943		298.12	0.6051
288.90	0.5857		291.40	0.5884		293.34	0.5915		294.99	0.5945		298.45	0.6122
289.01	0.5893		291.43	0.5885		293.46	0.5956		295.02	0.5982		298.55	0.6029
289.06	0.5876		291.47	0.5961		293.49	0.5925		295.14	0.6015		298.63	0.6053
289.11	0.5863		291.51	0.5935		293.55	0.5967		295.15	0.5963		298.82	0.6075
289.16	0.5857		291.56	0.5910		293.55	0.5971		295.19	0.5972		298.84	0.6072
289.18	0.5864		291.65	0.5892		293.58	0.5958		295.23	0.5953		299.13	0.6044
289.22	0.5894		291.68	0.5921		293.64	0.5935		295.26	0.6026		299.44	0.6038
289.29	0.5861		291.71	0.5902		293.65	0.5929		295.32	0.5983		299.49	0.6089
289.32	0.5871		291.77	0.5909		293.70	0.5941		295.36	0.5935		299.59	0.6074
289.43	0.5914		291.85	0.5938		293.71	0.5955		295.40	0.5960		299.62	0.6115
289.47	0.5866		291.87	0.5940		293.71	0.5972		295.45	0.5996		299.75	0.6096
289.51	0.5872		291.92	0.5946		293.84	0.5906		295.47	0.5967		300.05	0.6083
289.56	0.5940		291.99	0.5928		293.92	0.5964		295.47	0.5928		300.13	0.6076
289.61	0.5881		292.05	0.5905		293.94	0.5930		294.68	0.5983		300.20	0.6129
289.65	0.5886		292.08	0.5896		293.99	0.5930		294.35	0.5939		300.25	0.6089
289.68	0.5885		292.15	0.5937		294.02	0.5981		294.27	0.5918			
289.78	0.5897		292.16	0.5908		294.07	0.5935		294.19	0.5980			
289.84	0.5861		292.22	0.5858		294.11	0.5977		294.33	0.5931			
289.88	0.5889		292.25	0.5957		294.16	0.5918		294.32	0.5894			
289.95	0.5891		292.29	0.5909		294.18	0.5942		294.35	0.5920			
290.00	0.5946		292.37	0.5917		294.21	0.5988		294.52	0.5941			
290.07	0.5907		292.45	0.5902		294.24	0.5963		294.64	0.5925			
290.12	0.5902		292.48	0.5910		294.26	0.5906		294.67	0.5962			
290.16	0.5925		292.50	0.5911		294.30	0.5956		294.70	0.5917			
290.20	0.5874		292.59	0.5937		294.38	0.5955		294.77	0.5915			
290.28	0.5899		292.58	0.5932		294.43	0.5909		294.93	0.5929			
290.63	0.5922		292.63	0.5930		294.46	0.5924		295.01	0.5958			
290.67	0.5880		292.70	0.5916		294.47	0.5958		295.20	0.5935			
290.74	0.5875		292.75	0.5925		294.51	0.5986		296.08	0.5924			
290.79	0.5864		292.85	0.5962		294.54	0.5933		296.15	0.5955			
290.81	0.5937		292.88	0.5918		294.55	0.5990		296.34	0.5978			
290.85	0.5896		292.91	0.5924		294.64	0.5957		296.50	0.5973			

	Th	ermal Con	ductivit	y V	alues fo	r Water	' in	W/m.K	(F	Run	-3)	
Т(К)	k	Т (К)	k		Т (К)	k		Т (К)	k			
273.28	0.5662	282.33	0.5768		290.42	0.5925		296.39	0.5965			
273.15	0.5646	282.53	0.5750		290.53	0.5884		296.50	0.5988			
273.12	0.5631	282.93	0.5740		290.63	0.5881		296.60	0.5993			
273.12	0.5650	283.15	0.5791		290.82	0.5924		296.76	0.6034			
273.31	0.5570	283.31	0.5797		290.93	0.5940		296.79	0.5994			
273.51	0.5560	283.51	0.5760		291.06	0.5935		297.13	0.6001			
273.74	0.5605	283.67	0.5754		291.25	0.5926		297.37	0.6027			
273.98	0.5568	283.89	0.5826		291.32	0.5902		297.49	0.6092			
274.16	0.5622	284.06	0.5799		291.76	0.5975		297.59	0.6031			
274.41	0.5612	284.20	0.5787		291.89	0.5964		297.71	0.6013			
274.62	0.5539	284.43	0.5827		292.06	0.5950		297.80	0.6045			
275.03	0.5637	284.59	0.5766		292.18	0.5958		297.93	0.6056			
275.30	0.5592	284.76	0.5794		292.25	0.5926		298.12	0.6069			
275.68	0.5611	284.94	0.5789		292.39	0.5965		298.30	0.6010			
275.83	0.5661	285.11	0.5786		292.70	0.5939		298.40	0.6054			
276.42	0.5613	285.30	0.5787		292.80	0.5923		298.48	0.6036			
276.64	0.5579	285.65	0.5811		292.96	0.5967		298.55	0.6071			
276.74	0.5582	285.81	0.5822		293.25	0.5940		298.61	0.6056			
276.90	0.5574	285.96	0.5843		293.32	0.5950		298.82	0.6112			
277.04	0.5577	286.17	0.5833		293.39	0.5967		298.91	0.6063			
277.26	0.5599	286.55	0.5820		293.56	0.5931		298.96	0.6068			
277.89	0.5651	286.68	0.5815		293.74	0.5949		299.08	0.6131			
278.19	0.5619	286.79	0.5847		293.84	0.5978		299.19	0.6147			
278.40	0.5653	287.30	0.5829		293.96	0.5983		299.24	0.6030			
278.59	0.5652	287.48	0.5858		294.19	0.5941		299.35	0.6120			
279.08	0.5665	287.64	0.5801		294.36	0.5974		299.53	0.6075			
279.31	0.5668	287.80	0.5850		294.72	0.5941		299.58	0.6039			
279.53	0.5685	288.01	0.5914		294.85	0.5958		299.67	0.6099			
279.72	0.5663	288.12	0.5871		295.06	0.5959		299.79	0.6066			
280.11	0.5712	288.46	0.5880		295.15	0.5932		299.94	0.6099			
280.34	0.5715	288.60	0.5835		295.37	0.5953		300.01	0.6114			
280.55	0.5710	288.77	0.5842		295.50	0.5949		300.20	0.6152			
280.76	0.5740	288.92	0.5867		295.71	0.5990						
280.91	0.5709	289.34	0.5882		295.81	0.5935						
281.17	0.5715	289.65	0.5917		295.90	0.5947						
281.34	0.5691	289.80	0.5891		296.00	0.5983						
281.55	0.5737	289.97	0.5879		296.08	0.5948						
281.93	0.5731	290.11	0.5915		296.20	0.5982						
282.12	0.5772	290.28	0.5896		296.30	0.5958						

	Thermal	Conducti	vity Val	ues	for Eth	ylene G	lyc	ol in W,	/m.K	(Run-1)	
Т (К)	k	Т (К)	k		Т (К)	k		Т (К)	k	Т (К)	k
293.18	0.2447	322.17	0.2507		308.77	0.2500		302.43	0.2493	297.91	0.2492
293.23	0.2462	321.53	0.2509		308.56	0.2499		302.33	0.2493	297.86	0.2475
293.26	0.2479	321.22	0.2520		308.39	0.2471		302.07	0.2502	297.65	0.2490
293.35	0.2465	320.92	0.2517		308.18	0.2512		301.97	0.2494	297.58	0.2489
293.41	0.2457	320.31	0.2527		308.01	0.2490		301.87	0.2500	297.51	0.2465
293.65	0.2468	319.99	0.2534		307.82	0.2514		301.65	0.2484	297.44	0.2480
293.80	0.2467	319.38	0.2500		307.66	0.2486		301.51	0.2492	297.41	0.2488
337.66	0.2533	319.09	0.2488		307.48	0.2521		301.26	0.2471	297.36	0.2513
337.16	0.2529	318.77	0.2494		307.33	0.2506		301.13	0.2492	297.30	0.2507
336.63	0.2533	318.20	0.2524		307.13	0.2471		301.01	0.2480	297.27	0.2499
336.12	0.2532	317.92	0.2487		306.95	0.2506		300.95	0.2497	297.16	0.2495
335.65	0.2550	317.68	0.2535		306.81	0.2495		300.83	0.2486	297.08	0.2469
335.16	0.2519	317.39	0.2508		306.63	0.2500		300.74	0.2473	297.01	0.2493
334.71	0.2577	317.13	0.2523		306.46	0.2515		300.64	0.2505	296.97	0.2485
334.23	0.2577	316.85	0.2492		306.27	0.2516		300.52	0.2467	296.87	0.2477
332.42	0.2556	316.59	0.2519		306.12	0.2507		300.43	0.2478	296.82	0.2501
331.96	0.2567	316.31	0.2497		305.97	0.2484		300.40	0.2507	296.80	0.2483
331.52	0.2549	316.08	0.2488		305.82	0.2489		300.25	0.2487	296.72	0.2494
331.11	0.2537	315.54	0.2498		305.67	0.2486		300.12	0.2481	296.67	0.2480
330.69	0.2537	314.78	0.2504		305.38	0.2481		300.02	0.2497	296.60	0.2474
330.29	0.2516	314.52	0.2506		305.22	0.2477		299.94	0.2482	296.47	0.2503
329.86	0.2509	314.28	0.2511		305.09	0.2480		299.82	0.2499	296.40	0.2483
329.46	0.2557	314.03	0.2480		304.93	0.2484		299.73	0.2499	296.38	0.2483
329.04	0.2515	313.81	0.2497		304.77	0.2512		299.67	0.2488	296.36	0.2476
328.66	0.2540	313.57	0.2528		304.64	0.2502		299.57	0.2496	296.27	0.2490
328.24	0.2518	313.11	0.2478		304.49	0.2472		299.48	0.2504	296.16	0.2492
327.86	0.2518	312.89	0.2501		304.34	0.2481		299.28	0.2482	296.12	0.2498
327.50	0.2522	312.45	0.2497		304.22	0.2507		299.19	0.2484	296.07	0.2481
327.11	0.2509	312.24	0.2503		304.04	0.2495		299.06	0.2484	296.01	0.2487
326.34	0.2552	311.99	0.2505		303.90	0.2513		298.97	0.2489	295.91	0.2495
325.96	0.2502	311.76	0.2489		303.75	0.2501		298.89	0.2480	295.88	0.2502
325.26	0.2508	311.54	0.2483		303.46	0.2513		298.78	0.2487	260.82	0.2485
324.90	0.2526	311.31	0.2498		303.34	0.2490		298.73	0.2469	260.97	0.2418
324.56	0.2547	311.11	0.2513		303.21	0.2481		298.61	0.2501	261.15	0.2412
324.22	0.2528	310.74	0.2494		303.08	0.2494		298.50	0.2476	261.23	0.2418
323.90	0.2546	310.56	0.2495		302.93	0.2490		298.32	0.2493	261.38	0.2425
323.54	0.2491	310.32	0.2518		302.81	0.2496		298.09	0.2494	261.73	0.2405
322.85	0.2512	309.32	0.2533		302.69	0.2511		298.00	0.2490	261.92	0.2423
322.53	0.2524	309.12	0.2498		302.56	0.2523		297.98	0.2492	262.14	0.2418

Appendix B: Tabulated Thermal Conductivity Results for Ethylene Glycol

	Therma	l Co	onductiv	vity Val	ues	for Eth	ylene G	lyc	ol in W	/m.K	(Run-1)	
Т(К)	k		T (K)	k		Т(К)	k		Т(К)	k	Т (К)	k
262.36	0.2410		274.16	0.2420		282.74	0.2453		289.29	0.2454	294.79	0.2480
262.59	0.2402		274.35	0.2433		283.01	0.2455		289.43	0.2438	294.87	0.2475
262.85	0.2402		274.72	0.2450		283.13	0.2464		289.51	0.2460	294.93	0.2479
263.35	0.2433		274.92	0.2425		283.29	0.2451		289.62	0.2458	295.01	0.2464
263.57	0.2412		275.13	0.2441		283.55	0.2448		289.72	0.2472	295.14	0.2464
263.84	0.2413		275.28	0.2437		283.71	0.2453		289.96	0.2466	295.23	0.2474
264.09	0.2417		275.47	0.2435		284.00	0.2448		290.05	0.2447	295.27	0.2473
264.34	0.2407		275.88	0.2437		284.14	0.2446		290.16	0.2457	295.34	0.2486
264.61	0.2395		276.05	0.2455		284.29	0.2461		290.26	0.2448	295.43	0.2475
264.87	0.2399		276.22	0.2433		284.54	0.2451		290.42	0.2462	295.59	0.2487
265.16	0.2396		276.41	0.2433		284.68	0.2451		290.52	0.2451	295.78	0.2482
265.41	0.2401		276.61	0.2452		284.82	0.2460		290.61	0.2451		
266.12	0.2417		276.76	0.2437		284.99	0.2459		290.80	0.2444		
266.40	0.2416		276.97	0.2430		285.26	0.2441		290.93	0.2458		
266.88	0.2413		277.87	0.2427		285.38	0.2452		291.22	0.2462		
267.15	0.2406		278.05	0.2437		285.66	0.2461		291.42	0.2457		
267.40	0.2421		278.24	0.2442		285.91	0.2436		291.56	0.2474		
267.60	0.2442		278.38	0.2436		286.03	0.2451		291.65	0.2462		
267.87	0.2398		278.54	0.2443		286.14	0.2448		291.75	0.2463		
268.10	0.2413		278.88	0.2460		286.30	0.2455		292.15	0.2460		
268.33	0.2425		279.07	0.2455		286.43	0.2461		292.23	0.2452		
268.58	0.2426		279.22	0.2450		286.56	0.2456		292.34	0.2461		
268.82	0.2425		279.41	0.2454		286.94	0.2445		292.43	0.2459		
269.31	0.2418		279.57	0.2436		287.08	0.2460		292.62	0.2463		
269.54	0.2424		279.75	0.2443		287.20	0.2458		292.91	0.2467		
269.74	0.2414		279.91	0.2457		287.32	0.2463		293.15	0.2463		
270.22	0.2413		280.24	0.2438		287.43	0.2447		293.26	0.2450		
270.44	0.2415		280.39	0.2441		287.58	0.2456		293.36	0.2476		
270.88	0.2411		280.54	0.2448		287.79	0.2461		293.43	0.2469		
271.10	0.2418		280.71	0.2454		288.09	0.2457		293.50	0.2456		
271.32	0.2443		280.90	0.2462		288.20	0.2452		293.61	0.2465		
271.76	0.2420		281.05	0.2450		288.33	0.2455		293.76	0.2462		
271.97	0.2431		281.17	0.2459		288.46	0.2461		293.89	0.2465		
272.18	0.2420		281.48	0.2453		288.55	0.2462		294.02	0.2480		
272.38	0.2427		281.65	0.2457		288.67	0.2444		294.26	0.2461		
272.80	0.2438		281.82	0.2457		288.80	0.2440		294.35	0.2465		
273.40	0.2433		282.08	0.2444		288.92	0.2446		294.45	0.2463		
273.58	0.2427		282.23	0.2443		289.01	0.2463		294.58	0.2471		
273.80	0.2425		282.57	0.2455		289.13	0.2456		294.66	0.2476		

	Thermal	l Conducti	vity Val	ues	for Eth	ylene G	lyc	ol in W	/m.K	(Run-2)	
Т (К)	k	Т (К)	k		Т (К)	k		Т (К)	k	Т (К)	k
350.50	0.2486	321.23	0.2488		309.73	0.2477		301.99	0.2494	264.48	0.2417
349.33	0.2491	320.93	0.2487		309.50	0.2519		301.88	0.2499	264.76	0.2405
348.25	0.2493	320.62	0.2505		309.32	0.2487		301.58	0.2500	264.99	0.2419
347.17	0.2492	320.32	0.2546		309.16	0.2501		301.46	0.2497	265.28	0.2417
346.16	0.2488	320.06	0.2486		309.01	0.2483		301.40	0.2491	266.03	0.2408
345.17	0.2528	319.82	0.2482		308.80	0.2499		301.28	0.2476	266.31	0.2407
344.23	0.2503	319.56	0.2509		308.52	0.2478		301.15	0.2479	266.57	0.2415
343.32	0.2485	319.30	0.2478		308.38	0.2507		300.76	0.2480	266.82	0.2406
342.44	0.2491	319.01	0.2511		308.19	0.2507		300.63	0.2496	267.07	0.2401
341.65	0.2520	318.77	0.2477		308.01	0.2494		300.45	0.2485	267.30	0.2410
340.89	0.2535	317.71	0.2524		307.92	0.2493		300.36	0.2494	267.54	0.2410
340.12	0.2523	317.46	0.2476		307.81	0.2490		300.28	0.2491	267.77	0.2419
339.37	0.2506	317.21	0.2497		307.62	0.2489		299.95	0.2485	268.01	0.2427
338.69	0.2528	316.94	0.2523		307.26	0.2509		299.85	0.2484	268.29	0.2422
338.01	0.2539	316.64	0.2489		306.93	0.2494		299.75	0.2492	268.52	0.2413
337.34	0.2521	316.39	0.2520		306.76	0.2500		299.67	0.2477	268.76	0.2414
336.07	0.2552	315.99	0.2490		306.53	0.2487		299.57	0.2487	269.26	0.2423
335.47	0.2515	315.75	0.2499		306.21	0.2492		299.48	0.2504	270.15	0.2414
334.92	0.2540	315.51	0.2498		305.97	0.2499		299.41	0.2480	270.44	0.2416
334.36	0.2504	315.33	0.2502		305.78	0.2487		299.32	0.2482	270.63	0.2438
333.82	0.2519	315.11	0.2507		305.59	0.2497		299.27	0.2501	270.84	0.2421
333.30	0.2561	314.88	0.2516		305.48	0.2488		299.16	0.2501	271.46	0.2436
330.41	0.2501	314.10	0.2507		305.29	0.2489		298.95	0.2497	271.67	0.2417
329.98	0.2495	313.94	0.2488		305.09	0.2497		298.84	0.2506	271.91	0.2435
329.59	0.2521	313.77	0.2511		304.91	0.2510		298.73	0.2488	272.12	0.2423
329.16	0.2537	313.27	0.2495		304.78	0.2482		298.64	0.2497	272.33	0.2427
327.93	0.2525	313.09	0.2514		304.35	0.2500		298.56	0.2479	272.52	0.2419
327.15	0.2491	312.90	0.2494		304.07	0.2495		261.50	0.2424	272.72	0.2429
325.37	0.2496	312.66	0.2502		303.94	0.2516		261.67	0.2426	272.96	0.2434
325.02	0.2530	312.48	0.2493		303.81	0.2503		261.77	0.2425	273.15	0.2432
324.68	0.2484	312.20	0.2498		303.58	0.2492		261.98	0.2411	273.35	0.2436
324.32	0.2500	311.79	0.2489		303.44	0.2499		262.15	0.2412	273.54	0.2436
324.01	0.2483	311.62	0.2504		303.35	0.2481		262.35	0.2422	273.74	0.2441
323.71	0.2502	311.24	0.2496		303.04	0.2480		262.59	0.2430	273.94	0.2447
323.41	0.2502	311.05	0.2513		302.82	0.2472		263.00	0.2406	274.15	0.2443
323.05	0.2483	310.66	0.2484		302.69	0.2488		263.25	0.2399	274.33	0.2427
322.74	0.2528	310.47	0.2491		302.52	0.2498		263.52	0.2404	274.51	0.2438
322.12	0.2523	310.13	0.2496		302.37	0.2482		263.80	0.2401	274.69	0.2438
321.81	0.2484	309.94	0.2471		302.26	0.2498		264.24	0.2401	274.89	0.2418

	Therma	l Co	onducti	vity Val	ues	for Eth	ylene G	lyc	ol in W	/m.K	(Run-2)	
T (K)	k		T (K)	k		Т (К)	k		T (K)	k	T (K)	k
275.06	0.2440		284.18	0.2438		290.92	0.2456		295.67	0.2472	299.70	0.2487
275.27	0.2434		284.31	0.2442		291.04	0.2446		295.74	0.2484	299.86	0.2496
275.46	0.2422		284.43	0.2454		291.29	0.2457		295.87	0.2482	299.91	0.2500
276.02	0.2448		284.57	0.2452		291.38	0.2463		295.96	0.2489	300.01	0.2487
276.19	0.2448		284.68	0.2447		291.50	0.2451		296.04	0.2471	300.03	0.2505
276.59	0.2441		284.83	0.2456		291.71	0.2458		296.11	0.2492	300.39	0.2488
276.76	0.2434		285.15	0.2444		291.82	0.2451		296.19	0.2498	300.46	0.2499
277.11	0.2437		285.28	0.2459		291.95	0.2461		296.28	0.2504	300.91	0.2500
277.28	0.2427		285.38	0.2460		292.06	0.2471		296.42	0.2473	301.15	0.2491
277.41	0.2450		285.52	0.2455		292.19	0.2454		296.48	0.2493	301.42	0.2499
277.60	0.2427		285.66	0.2463		292.38	0.2466		296.60	0.2481	301.50	0.2488
278.13	0.2447		285.79	0.2462		292.45	0.2476		296.75	0.2497	301.55	0.2493
278.46	0.2446		286.06	0.2455		292.53	0.2463		296.87	0.2492	301.67	0.2496
278.65	0.2435		286.19	0.2452		292.71	0.2476		296.94	0.2468	301.90	0.2489
278.82	0.2431		286.46	0.2465		292.82	0.2455		297.13	0.2485	302.13	0.2500
279.19	0.2435		286.97	0.2463		293.13	0.2472		297.22	0.2494	302.26	0.2506
279.31	0.2444		287.11	0.2441		293.19	0.2483		297.28	0.2490	302.35	0.2491
279.48	0.2431		287.21	0.2449		293.39	0.2464		297.35	0.2486	302.92	0.2505
279.64	0.2450		287.35	0.2447		293.62	0.2470		297.52	0.2475	303.07	0.2494
279.82	0.2438		287.57	0.2456		293.71	0.2474		297.61	0.2465		
280.30	0.2445		287.67	0.2453		293.78	0.2460		297.69	0.2492		
280.46	0.2439		287.79	0.2453		293.90	0.2473		297.77	0.2480		
280.60	0.2436		287.92	0.2450		294.00	0.2475		297.82	0.2467		
280.77	0.2453		288.05	0.2459		294.17	0.2479		297.85	0.2490		
281.06	0.2443		288.39	0.2451		294.35	0.2488		297.96	0.2497		
281.20	0.2448		288.53	0.2466		294.43	0.2483		298.02	0.2503		
281.53	0.2447		288.61	0.2467		294.49	0.2466		298.22	0.2483		
281.67	0.2445		288.76	0.2449		294.55	0.2482		298.31	0.2487		
281.84	0.2443		289.19	0.2463		294.62	0.2488		298.38	0.2508		
282.00	0.2437		289.44	0.2459		294.73	0.2465		298.46	0.2503		
282.14	0.2442		289.69	0.2454		294.79	0.2475		298.54	0.2487		
282.42	0.2464		289.75	0.2452		294.87	0.2486		298.61	0.2504		
282.77	0.2440		289.84	0.2458		294.96	0.2485		298.91	0.2488		
282.89	0.2460		290.13	0.2454		295.06	0.2476		298.99	0.2483		
283.15	0.2451		290.24	0.2456		295.14	0.2487		299.04	0.2499		
283.48	0.2445		290.43	0.2457		295.28	0.2474		299.04	0.2495		
283.58	0.2445		290.62	0.2460		295.36	0.2478		299.07	0.2499		
283.89	0.2458		290.75	0.2453		295.45	0.2484		299.34	0.2495		
284.02	0.2441		290.83	0.2465		295.59	0.2469		299.47	0.2480		

	Thermal	l Conducti	vity Val	ues	for Eth	ylene G	lyc	ol in W	/m.K	(Run-3)	
Т (К)	k	Т (К)	k		Т (К)	k		Т (К)	k	Т (К)	k
293.31	0.2444	325.66	0.2505		309.77	0.2508		302.95	0.2502	297.70	0.2477
293.28	0.2487	324.73	0.2547		309.43	0.2492		302.83	0.2491	297.54	0.2489
293.27	0.2476	324.42	0.2546		308.81	0.2506		302.73	0.2509	297.44	0.2498
293.28	0.2470	324.12	0.2500		308.35	0.2509		302.62	0.2506	297.36	0.2504
293.29	0.2464	323.81	0.2543		307.90	0.2475		302.47	0.2493	297.26	0.2484
293.29	0.2470	322.55	0.2535		307.72	0.2486		302.37	0.2480	297.21	0.2480
293.29	0.2482	321.87	0.2515		307.53	0.2489		302.23	0.2479	297.15	0.2501
293.29	0.2475	321.38	0.2495		307.22	0.2504		302.10	0.2474	297.07	0.2479
293.29	0.2462	321.09	0.2509		307.07	0.2512		301.99	0.2508	297.02	0.2497
293.27	0.2479	320.19	0.2512		306.78	0.2476		301.87	0.2502	296.91	0.2497
293.26	0.2474	319.88	0.2516		306.61	0.2505		301.36	0.2485	296.85	0.2492
293.23	0.2479	319.33	0.2501		306.48	0.2500		301.30	0.2479	296.79	0.2481
293.22	0.2484	318.81	0.2503		305.99	0.2478		301.20	0.2502	296.67	0.2500
293.23	0.2469	318.60	0.2508		309.46	0.2488		301.07	0.2486	296.58	0.2471
293.24	0.2487	318.28	0.2508		309.12	0.2506		300.93	0.2503	296.46	0.2497
293.24	0.2479	317.98	0.2519		308.94	0.2509		300.69	0.2493	296.26	0.2495
293.25	0.2462	317.78	0.2497		308.76	0.2493		300.54	0.2502	296.13	0.2487
293.24	0.2468	317.52	0.2511		308.56	0.2500		300.25	0.2485	296.02	0.2482
293.24	0.2470	317.29	0.2525		308.17	0.2489		300.12	0.2500	260.59	0.2415
293.24	0.2487	316.83	0.2527		307.81	0.2485		300.04	0.2474	260.68	0.2410
293.23	0.2486	316.56	0.2490		307.64	0.2468		299.95	0.2509	260.82	0.2430
338.10	0.2525	316.30	0.2487		307.46	0.2501		299.73	0.2501	261.08	0.2415
336.23	0.2536	316.03	0.2509		307.15	0.2517		299.60	0.2501	261.25	0.2418
335.32	0.2556	315.83	0.2497		306.96	0.2513		299.30	0.2508	261.45	0.2405
334.86	0.2541	315.63	0.2497		306.09	0.2510		299.21	0.2494	261.62	0.2415
334.40	0.2551	314.71	0.2518		305.95	0.2492		299.09	0.2493	262.06	0.2415
333.54	0.2518	314.04	0.2487		305.75	0.2498		298.94	0.2503	262.28	0.2404
332.30	0.2522	313.86	0.2516		305.42	0.2512		298.72	0.2485	262.50	0.2413
331.89	0.2545	313.63	0.2484		305.30	0.2496		298.64	0.2497	262.76	0.2425
331.50	0.2541	313.15	0.2478		304.97	0.2504		298.51	0.2503	262.98	0.2419
331.10	0.2552	312.70	0.2505		304.83	0.2499		298.44	0.2469	263.22	0.2415
330.69	0.2520	312.49	0.2485		304.24	0.2503		298.39	0.2468	263.45	0.2405
329.21	0.2550	312.30	0.2492		304.05	0.2495		298.34	0.2500	263.73	0.2402
328.82	0.2520	312.03	0.2500		303.92	0.2512		298.26	0.2498	263.96	0.2402
328.46	0.2523	311.89	0.2507		303.80	0.2487		298.07	0.2483	264.19	0.2402
328.13	0.2533	311.67	0.2509		303.66	0.2512		298.00	0.2490	264.44	0.2409
327.44	0.2532	311.12	0.2499		303.38	0.2511		297.93	0.2504	264.68	0.2430
327.09	0.2503	310.33	0.2492		303.24	0.2497		297.86	0.2497	264.92	0.2400
326.36	0.2531	310.15	0.2496		303.09	0.2483		297.80	0.2480	265.22	0.2412

	Therma	l Co	onducti	vity Val	ues	for Eth	ylene G	lyc	ol in W	/m.K	(Run-3)	
Т (К)	k		Т (К)	k		Т (К)	k		Т (К)	k		Т (К)	k
265.50	0.2423		275.37	0.2436		283.67	0.2464		289.47	0.2458		294.60	0.2470
265.75	0.2416		275.54	0.2439		283.83	0.2433		289.58	0.2461		294.76	0.2466
265.99	0.2423		275.71	0.2425		283.95	0.2449		289.68	0.2460		294.83	0.2478
266.22	0.2415		275.93	0.2428		284.09	0.2461		289.88	0.2459		294.89	0.2470
266.45	0.2426		276.08	0.2452		284.25	0.2445		289.99	0.2456		294.99	0.2484
266.69	0.2423		276.26	0.2425		284.38	0.2453		290.12	0.2450		295.04	0.2479
266.95	0.2412		276.45	0.2428		284.52	0.2456		290.23	0.2459		295.14	0.2467
267.17	0.2402		276.64	0.2427		284.93	0.2446		290.32	0.2450		295.29	0.2480
267.39	0.2412		277.03	0.2429		285.07	0.2458		290.65	0.2447		295.42	0.2488
267.62	0.2402		277.16	0.2430		285.20	0.2444		290.75	0.2451		295.49	0.2492
267.88	0.2412		277.37	0.2453		285.33	0.2451		291.00	0.2460		295.56	0.2475
268.12	0.2412		277.68	0.2445		285.48	0.2464		291.09	0.2460		295.62	0.2491
268.38	0.2415		277.88	0.2427		285.74	0.2452		291.21	0.2468		295.76	0.2477
268.59	0.2429		278.01	0.2451		285.88	0.2444		291.30	0.2456		295.94	0.2485
269.02	0.2423		278.17	0.2453		286.00	0.2459		291.40	0.2469		296.08	0.2472
269.27	0.2423		278.36	0.2446		286.14	0.2452		291.51	0.2452		296.12	0.2481
269.49	0.2410		278.51	0.2443		286.28	0.2450		291.79	0.2457		296.22	0.2480
269.71	0.2421		278.68	0.2434		286.52	0.2457		292.00	0.2458		296.44	0.2466
270.42	0.2423		279.03	0.2442		286.76	0.2450		292.10	0.2472		296.54	0.2493
270.60	0.2433		279.22	0.2451		286.87	0.2459		292.26	0.2474		296.63	0.2475
270.82	0.2434		279.56	0.2436		287.02	0.2460		292.51	0.2452		296.71	0.2479
271.03	0.2423		279.78	0.2444		287.14	0.2461		292.70	0.2473		296.78	0.2475
271.26	0.2429		280.24	0.2442		287.27	0.2450		292.81	0.2475		296.84	0.2497
271.86	0.2424		280.57	0.2446		287.41	0.2444		293.03	0.2453		296.92	0.2488
272.05	0.2441		280.83	0.2453		287.55	0.2461		293.11	0.2462		296.99	0.2489
272.26	0.2441		281.18	0.2442		287.66	0.2457		293.20	0.2468		297.05	0.2475
272.66	0.2439		281.33	0.2441		287.77	0.2448		293.35	0.2461		297.15	0.2480
272.90	0.2444		281.51	0.2442		287.89	0.2461		293.59	0.2469		297.22	0.2473
273.08	0.2443		281.64	0.2443		288.02	0.2449		293.65	0.2459		297.30	0.2494
273.29	0.2437		281.78	0.2449		288.16	0.2455		293.76	0.2471		297.36	0.2501
273.47	0.2438		281.94	0.2441		288.28	0.2458		293.84	0.2474		297.42	0.2470
273.68	0.2411		282.09	0.2445		288.40	0.2468		293.89	0.2468		297.50	0.2490
273.87	0.2431		282.36	0.2441		288.50	0.2456		293.96	0.2459		297.58	0.2482
274.11	0.2423		282.50	0.2453		288.63	0.2459		294.08	0.2468		297.64	0.2490
274.28	0.2430		282.66	0.2449		288.74	0.2449		294.17	0.2466		297.73	0.2479
274.67	0.2439		282.81	0.2440		288.87	0.2458		294.26	0.2477		297.81	0.2491
274.82	0.2435		282.97	0.2453		288.99	0.2448		294.36	0.2473		297.87	0.2502
275.02	0.2436		283.12	0.2458		289.13	0.2458		294.45	0.2469		298.09	0.2503
275.20	0.2447		283.26	0.2462		289.27	0.2456		294.53	0.2482		298.43	0.2500

	Therma	l Co	onducti	vity Val	ues	for Eth	ylene G	lyc	ol in W	/m.K		(Run-3)	
Т (К)	k												
298.69	0.2492												
298.74	0.2501												
298.90	0.2491												
298.97	0.2502												
299.40	0.2493												
299.93	0.2509												
300.56	0.2505												
300.69	0.2483												
300.96	0.2498												
301.36	0.2497												
301.40	0.2500												
301.98	0.2494												
302.46	0.2507												
302.83	0.2500												
303.51	0.2509												
303.76	0.2501												
303.77	0.2500												
		_											
		_											
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	Ther	mal Condu	uctivity	Va	lues for	Glycero	ol ir	n W/m.H	Κ ([Ru	n-1)	
Т (К)	k	Т (К)	k		Т (К)	k		T (K)	k		Т (К)	k
294.19	0.2755	304.28	0.2769		310.77	0.2770		316.09	0.2778		319.57	0.2790
294.14	0.2760	304.54	0.2764		310.81	0.2765		316.19	0.2773		319.54	0.2787
294.23	0.2752	304.82	0.2748		310.86	0.2767		316.29	0.2778		319.64	0.2770
294.29	0.2753	305.36	0.2764		310.91	0.2786		316.39	0.2782		319.74	0.2779
294.60	0.2767	305.68	0.2741		311.03	0.2786		316.46	0.2785		319.67	0.2787
294.78	0.2755	305.90	0.2768		311.11	0.2763		316.61	0.2754		319.62	0.2771
294.98	0.2753	306.18	0.2742		311.30	0.2769		316.73	0.2786		319.63	0.2789
295.25	0.2765	306.44	0.2765		311.47	0.2772		316.78	0.2764		319.62	0.2790
295.50	0.2731	306.66	0.2744		311.62	0.2792		316.83	0.2775		319.66	0.2774
295.76	0.2737	306.84	0.2749		311.78	0.2779		316.96	0.2777		319.61	0.2776
295.98	0.2765	307.01	0.2751		311.88	0.2763		317.07	0.2781		319.63	0.2787
296.23	0.2762	307.29	0.2748		312.01	0.2769		317.17	0.2782		319.60	0.2781
296.47	0.2728	307.50	0.2761		312.10	0.2773		317.30	0.2772		319.71	0.2776
296.66	0.2757	307.70	0.2761		312.30	0.2772		317.32	0.2770		319.71	0.2780
296.90	0.2762	307.84	0.2778		312.75	0.2765		317.48	0.2770		319.75	0.2783
297.10	0.2755	307.97	0.2763		312.89	0.2771		317.69	0.2782		319.73	0.2794
297.35	0.2763	308.15	0.2778		312.96	0.2765		317.77	0.2795		319.79	0.2781
297.57	0.2733	308.27	0.2768		313.05	0.2786		318.01	0.2782		319.76	0.2782
297.74	0.2758	308.39	0.2777		313.56	0.2768		318.12	0.2771		319.74	0.2781
297.95	0.2765	308.47	0.2782		313.69	0.2778		318.23	0.2771		319.79	0.2795
298.20	0.2750	308.66	0.2746		313.78	0.2763		318.34	0.2772		319.77	0.2794
298.44	0.2742	308.72	0.2765		313.98	0.2764		318.44	0.2778		319.85	0.2789
298.68	0.2757	308.73	0.2771		314.15	0.2778		318.57	0.2779		319.83	0.2788
298.87	0.2746	308.84	0.2769		314.30	0.2770		318.67	0.2788		319.88	0.2789
299.43	0.2760	308.88	0.2779		314.36	0.2769		318.77	0.2776		319.91	0.2786
299.65	0.2762	308.96	0.2788		314.59	0.2773		318.85	0.2761		319.87	0.2780
299.90	0.2750	309.06	0.2770		314.69	0.2778		318.93	0.2794		319.89	0.2781
300.16	0.2755	309.20	0.2786		314.69	0.2768		318.92	0.2785		319.85	0.2793
300.38	0.2739	309.30	0.2769		314.90	0.2789		319.06	0.2772		319.88	0.2773
300.64	0.2752	309.57	0.2772		314.94	0.2764		319.18	0.2775		320.04	0.2772
300.97	0.2752	309.70	0.2767		315.04	0.2779		319.22	0.2773		320.25	0.2774
301.29	0.2749	309.83	0.2768		315.12	0.2771		319.24	0.2782		320.48	0.2790
301.84	0.2742	310.04	0.2769		315.24	0.2760		319.26	0.2791		320.58	0.2778
302.05	0.2744	310.11	0.2766		315.29	0.2783		319.33	0.2757		320.67	0.2793
302.31	0.2743	310.23	0.2773		315.37	0.2761		319.36	0.2773		320.78	0.2785
302.64	0.2735	310.34	0.2778		315.41	0.2780		319.41	0.2769		320.79	0.2767
302.93	0.2746	310.44	0.2767		315.72	0.2763		319.45	0.2780		320.97	0.2772
303.43	0.2728	310.53	0.2764		315.75	0.2763		319.44	0.2770		321.06	0.2789
303.93	0.2770	310.64	0.2757		315.99	0.2763		319.44	0.2790		321.12	0.2788

Appendix C: Tabulated Thermal Conductivity Results for Glycerol

	The	rm	al Cond	uctivity	Va	lues for	Glycero	ol iı	n W/m.l	К (Ru	n-1)	
T (K)	k		Т (К)	k		Т (К)	k		Т (К)	k		Т (К)	k
321.17	0.2773		322.27	0.2770		330.52	0.2792		335.58	0.2786		344.39	0.2811
321.28	0.2778		322.63	0.2784		330.58	0.2801		335.73	0.2796		344.60	0.2817
321.39	0.2776		322.93	0.2780		330.63	0.2779		335.89	0.2781		344.68	0.2816
321.54	0.2770		323.21	0.2784		330.66	0.2796		336.07	0.2781		344.91	0.2792
321.81	0.2776		323.78	0.2773		330.74	0.2782		336.22	0.2797		345.15	0.2808
321.97	0.2797		324.03	0.2792		330.80	0.2801		337.04	0.2803		345.38	0.2819
322.08	0.2792		324.73	0.2795		330.85	0.2784		337.17	0.2825		345.27	0.2783
322.17	0.2795		324.97	0.2782		330.88	0.2804		337.40	0.2806		345.72	0.2811
322.60	0.2778		325.23	0.2780		330.91	0.2812		337.73	0.2791		345.65	0.2815
323.11	0.2764		325.58	0.2793		330.92	0.2792		337.95	0.2802		345.78	0.2792
323.68	0.2768		325.84	0.2791		331.03	0.2826		338.05	0.2806		345.80	0.2789
324.21	0.2774		325.96	0.2784		331.13	0.2799		338.13	0.2796		345.91	0.2829
324.44	0.2792		326.09	0.2804		331.22	0.2803		338.22	0.2790		346.44	0.2815
324.61	0.2785		326.20	0.2797		331.35	0.2815		338.34	0.2786		346.72	0.2823
324.86	0.2781		326.56	0.2791		331.54	0.2810		338.65	0.2784		347.23	0.2794
325.09	0.2773		326.72	0.2795		331.92	0.2810		338.90	0.2807		347.36	0.2801
325.29	0.2787		326.85	0.2795		331.98	0.2792		339.03	0.2785		347.58	0.2789
325.51	0.2787		326.97	0.2807		332.14	0.2787		339.18	0.2791		347.69	0.2822
325.70	0.2779		327.11	0.2789		332.31	0.2801		340.16	0.2788		347.72	0.2805
326.17	0.2780		327.37	0.2800		332.44	0.2776		340.73	0.2781		347.95	0.2789
326.32	0.2786		327.50	0.2785		332.56	0.2778		340.92	0.2797		348.00	0.2827
326.45	0.2786		327.83	0.2791		332.72	0.2779		341.01	0.2797		348.22	0.2806
326.61	0.2783		328.40	0.2797		332.84	0.2807		341.29	0.2800		348.38	0.2794
326.79	0.2786		328.70	0.2804		333.02	0.2803		341.64	0.2789		348.63	0.2807
326.87	0.2780		328.73	0.2776		333.12	0.2815		341.69	0.2784		349.09	0.2794
327.06	0.2789		329.05	0.2785		333.23	0.2799		341.80	0.2800			
327.27	0.2801		329.07	0.2803		333.37	0.2784		342.04	0.2823			
327.40	0.2782		329.19	0.2786		333.51	0.2799		342.24	0.2781			
327.56	0.2778		329.27	0.2811		333.71	0.2783		342.33	0.2800			
327.64	0.2785		329.45	0.2785		333.79	0.2780		342.39	0.2805			
327.82	0.2784		329.48	0.2788		333.94	0.2812		342.60	0.2788			
327.91	0.2783		329.78	0.2802		333.99	0.2793		342.73	0.2812			
328.35	0.2808		329.92	0.2798		334.18	0.2788		343.02	0.2817			
328.64	0.2793		330.01	0.2800		334.28	0.2787		343.10	0.2788			
328.81	0.2782		330.14	0.2802		334.43	0.2786		343.60	0.2813			
328.66	0.2790		330.17	0.2801		334.62	0.2800		343.69	0.2780			
320.61	0.2767		330.31	0.2794		334.81	0.2782		343.67	0.2826			
321.09	0.2759		330.39	0.2794		335.02	0.2801		343.90	0.2783			
321.44	0.2759		330.47	0.2781		335.35	0.2789		343.97	0.2793			

	Ther	rmal Cond	uctivity	Va	lues for	Glycero	ol ir	n W/m.I	К (Ru	n-2)	
T (K)	k	Т (К)	k		Т (К)	k		Т (К)	k		Т (К)	k
292.10	0.2764	301.10	0.2757		309.73	0.2769		315.13	0.2766		323.55	0.2786
292.11	0.2753	301.22	0.2745		309.83	0.2765		315.22	0.2786		323.90	0.2766
292.19	0.2774	301.34	0.2768		309.95	0.2767		315.40	0.2781		324.12	0.2780
292.36	0.2762	301.48	0.2761		310.06	0.2769		315.42	0.2762		324.47	0.2787
292.48	0.2763	301.62	0.2768		310.12	0.2751		315.53	0.2780		324.76	0.2772
292.90	0.2750	301.78	0.2761		310.39	0.2774		315.64	0.2759		325.18	0.2776
293.12	0.2753	302.08	0.2770		310.45	0.2762		315.87	0.2770		325.58	0.2766
293.37	0.2747	302.20	0.2761		310.63	0.2773		315.98	0.2775		326.85	0.2777
293.64	0.2748	302.43	0.2755		310.78	0.2763		316.14	0.2780		327.52	0.2780
293.83	0.2763	302.62	0.2757		310.97	0.2772		316.35	0.2758		327.78	0.2781
294.17	0.2769	302.75	0.2753		311.27	0.2764		316.52	0.2769		327.98	0.2801
294.32	0.2757	302.94	0.2739		311.57	0.2778		316.93	0.2765		328.28	0.2775
294.45	0.2751	303.14	0.2770		311.94	0.2767		317.02	0.2774		328.45	0.2789
294.58	0.2763	303.41	0.2757		312.16	0.2769		317.54	0.2754		328.63	0.2786
294.68	0.2770	303.77	0.2757		312.30	0.2751		317.81	0.2757		328.78	0.2791
294.80	0.2763	303.92	0.2771		312.53	0.2755		318.05	0.2779		328.96	0.2783
294.92	0.2780	304.42	0.2777		312.94	0.2793		318.48	0.2762		329.11	0.2798
295.48	0.2780	305.01	0.2756		313.19	0.2789		318.70	0.2786		329.18	0.2784
295.81	0.2766	305.20	0.2740		313.41	0.2787		319.09	0.2773		329.26	0.2781
295.98	0.2777	305.64	0.2749		313.49	0.2773		319.21	0.2768		329.40	0.2801
296.13	0.2752	306.08	0.2752		313.58	0.2762		319.47	0.2774		329.42	0.2791
296.30	0.2751	306.36	0.2749		313.62	0.2777		319.60	0.2786		329.77	0.2807
296.55	0.2743	306.63	0.2765		313.73	0.2785		319.81	0.2771		329.98	0.2809
296.81	0.2764	306.79	0.2764		313.90	0.2766		319.91	0.2783		330.05	0.2784
297.04	0.2766	306.94	0.2766		313.99	0.2779		320.11	0.2794		330.14	0.2793
297.28	0.2765	307.20	0.2775		314.06	0.2775		320.29	0.2792		330.13	0.2789
297.71	0.2751	307.35	0.2760		314.14	0.2769		320.24	0.2785		330.26	0.2795
297.88	0.2775	307.53	0.2771		314.24	0.2770		320.39	0.2795		330.30	0.2803
298.09	0.2740	307.80	0.2765		314.40	0.2774		320.49	0.2782		330.32	0.2782
298.29	0.2750	308.01	0.2767		314.51	0.2777		320.57	0.2779		330.39	0.2797
298.72	0.2777	308.14	0.2778		314.56	0.2771		320.67	0.2790		330.41	0.2790
298.94	0.2748	308.32	0.2780		314.65	0.2796		321.00	0.2785		330.49	0.2801
299.16	0.2749	308.63	0.2769		314.71	0.2785		321.31	0.2775		330.54	0.2805
299.38	0.2773	308.77	0.2772		314.79	0.2777		321.61	0.2765		330.55	0.2786
299.56	0.2749	308.98	0.2774		314.89	0.2768		321.93	0.2778		330.59	0.2790
300.01	0.2752	309.11	0.2755		314.91	0.2784		322.35	0.2761		331.45	0.2801
300.29	0.2770	309.23	0.2764		315.01	0.2773		322.59	0.2771		331.61	0.2790
300.69	0.2760	309.43	0.2762		315.07	0.2770		323.00	0.2762		331.76	0.2791
300.88	0.2767	309.57	0.2763		315.11	0.2789		323.30	0.2763		331.96	0.2798

	The	rm	al Condi	uctivity	Va	lues for	Glycero	ol iı	n W/m.l	К (Ru	n-3)	
T (K)	k		T (K)	k		Т (К)	k		Т (К)	k		Т (К)	k
292.88	0.2727		313.19	0.2786		324.22	0.2785		334.03	0.2781		341.91	0.2819
293.25	0.2755		313.41	0.2778		324.42	0.2820		334.15	0.2804		341.99	0.2809
293.47	0.2741		313.66	0.2795		324.72	0.2804		334.30	0.2791		342.03	0.2835
294.47	0.2739		313.92	0.2771		325.25	0.2811		334.39	0.2821		342.13	0.2804
294.85	0.2741		314.16	0.2790		325.67	0.2814		334.54	0.2802		342.24	0.2807
295.18	0.2740		314.42	0.2768		325.97	0.2808		334.90	0.2796		342.36	0.2833
295.56	0.2757		314.67	0.2782		326.21	0.2783		335.09	0.2798		342.35	0.2812
296.00	0.2732		314.93	0.2790		326.48	0.2807		335.32	0.2778		342.41	0.2804
296.49	0.2766		315.45	0.2792		326.68	0.2810		335.55	0.2805		342.44	0.2823
296.89	0.2744		315.68	0.2792		326.92	0.2811		335.74	0.2799		342.44	0.2803
297.29	0.2723		315.90	0.2775		327.15	0.2785		336.01	0.2808		342.66	0.2824
297.79	0.2757		316.17	0.2782		327.61	0.2824		336.08	0.2813		342.70	0.2802
298.29	0.2719		316.41	0.2790		327.82	0.2782		336.32	0.2829		342.69	0.2808
298.69	0.2736		316.59	0.2783		328.22	0.2778		336.51	0.2796		342.89	0.2818
299.13	0.2722		316.81	0.2799		329.05	0.2797		336.84	0.2800		342.87	0.2815
299.45	0.2754		317.23	0.2796		329.21	0.2785		337.18	0.2813		342.91	0.2835
300.84	0.2736		317.40	0.2795		329.55	0.2815		337.44	0.2782		342.97	0.2774
302.57	0.2760		317.61	0.2804		329.73	0.2796		337.77	0.2827		343.19	0.2802
302.99	0.2732		318.01	0.2788		329.90	0.2810		338.12	0.2833		343.23	0.2826
303.46	0.2750		318.21	0.2794		330.09	0.2798		338.52	0.2811		343.24	0.2817
304.65	0.2764		318.62	0.2795		330.25	0.2807		338.56	0.2780		343.27	0.2815
305.02	0.2742		318.81	0.2776		330.38	0.2788		338.77	0.2813		343.25	0.2830
305.36	0.2764		319.21	0.2800		330.57	0.2827		338.94	0.2829		343.21	0.2784
305.76	0.2747		319.38	0.2792		330.71	0.2796		339.05	0.2830		343.34	0.2808
306.15	0.2748		319.59	0.2778		330.84	0.2800		339.18	0.2814		343.40	0.2837
306.88	0.2742		319.72	0.2806		331.04	0.2791		339.44	0.2801		343.36	0.2839
307.13	0.2766		319.92	0.2788		331.33	0.2795		339.64	0.2817		343.57	0.2790
307.51	0.2757		320.15	0.2799		331.49	0.2778		339.94	0.2814		343.63	0.2825
307.89	0.2766		320.31	0.2808		331.61	0.2790		340.07	0.2778		343.73	0.2809
308.19	0.2766		320.90	0.2802		331.83	0.2781		340.26	0.2823		343.74	0.2798
308.54	0.2773		321.09	0.2810		332.00	0.2827		340.51	0.2808		343.75	0.2819
308.89	0.2769		321.27	0.2798		332.14	0.2802		340.62	0.2796		344.00	0.2806
310.24	0.2767		321.46	0.2800		332.29	0.2792		340.83	0.2817		344.00	0.2835
310.85	0.2761		321.91	0.2793		332.46	0.2812		340.98	0.2831		344.20	0.2795
311.41	0.2774		322.18	0.2801		332.60	0.2802		341.04	0.2812		344.23	0.2819
311.73	0.2779		322.45	0.2796		332.85	0.2782		341.22	0.2791		344.45	0.2806
312.04	0.2760		322.61	0.2793		333.13	0.2791		341.50	0.2784		344.63	0.2828
312.36	0.2766		323.14	0.2820		333.50	0.2787		341.58	0.2822		344.72	0.2829
312.87	0.2768		323.71	0.2791		333.83	0.2823		341.65	0.2814		344.95	0.2820

	The	rm	al Cond	uctivity	Va	lues for	Glycero	ol ir	n W/m.l	К (Ru	n-3)	
Т (К)	k		Т (К)	k		Т (К)	k		Т (К)	k		T (K)	k
345.13	0.2779												
345.18	0.2793												
345.36	0.2783												
345.70	0.2768												
345.87	0.2795												
					_								
					-						-		
					-								
					-								
					-								
					-						-		
											-		

Г	hermal	Со	nductiv	ity Valu	ies	for Pro	pylene (Glyo	col in W	/m.K	(Run-1)
T (K)	k		Т (К)	k		Т (К)	k		Т (К)	k	T (K)	k
238.88	0.1855		270.06	0.1875		287.35	0.1895		296.48	0.1899	293.73	0.1905
239.00	0.1848		270.91	0.1879		287.66	0.1895		296.64	0.1912	293.87	0.1888
239.19	0.1837		271.31	0.1887		287.84	0.1882		296.94	0.1904	294.09	0.1889
239.49	0.1840		271.72	0.1882		288.54	0.1896		297.21	0.1901	294.25	0.1891
239.78	0.1828		272.52	0.1882		288.79	0.1900		297.44	0.1919	294.39	0.1904
240.20	0.1820		272.91	0.1887		289.02	0.1882		297.60	0.1912	294.51	0.1906
240.75	0.1821		273.31	0.1874		289.71	0.1883		297.87	0.1913	294.64	0.1908
241.23	0.1821		274.11	0.1879		290.13	0.1892		298.02	0.1913	294.72	0.1902
241.72	0.1823		274.47	0.1871		288.09	0.1898		298.43	0.1911	294.79	0.1908
242.27	0.1820		274.81	0.1876		288.02	0.1899		298.55	0.1906	294.85	0.1909
250.92	0.1846		275.54	0.1884		288.07	0.1896		298.66	0.1914	294.86	0.1889
251.59	0.1835		275.93	0.1891		288.20	0.1900		298.78	0.1920	294.95	0.1906
252.82	0.1830		276.30	0.1880		288.28	0.1900		298.90	0.1916	294.99	0.1908
253.47	0.1847		277.68	0.1877		288.83	0.1893		299.17	0.1912	295.00	0.1910
254.15	0.1852		278.02	0.1893		289.19	0.1904		299.29	0.1907	295.07	0.1905
254.73	0.1850		278.37	0.1881		289.43	0.1900		299.39	0.1911	295.07	0.1911
255.35	0.1849		278.72	0.1899		289.66	0.1886		299.52	0.1911	295.03	0.1913
256.52	0.1853		279.03	0.1888		290.09	0.1901		299.63	0.1916	295.02	0.1908
257.09	0.1853		280.01	0.1883		290.29	0.1902		299.74	0.1923	294.96	0.1914
257.63	0.1857		280.31	0.1898		290.53	0.1897		300.00	0.1912	294.90	0.1913
258.19	0.1852		280.67	0.1887		290.72	0.1891		300.10	0.1912	294.88	0.1900
258.75	0.1851		280.91	0.1893		291.18	0.1899		300.22	0.1906	294.89	0.1907
259.27	0.1862		281.26	0.1881		291.99	0.1884		300.31	0.1915	294.85	0.1903
259.85	0.1856		281.89	0.1893		292.40	0.1887		300.41	0.1914	294.91	0.1905
260.37	0.1863		282.16	0.1898		292.93	0.1889		300.51	0.1917	294.97	0.1911
260.87	0.1862		282.77	0.1885		293.14	0.1893		300.95	0.1905	295.06	0.1910
261.42	0.1852		283.05	0.1883		293.28	0.1893		301.02	0.1907	295.10	0.1904
261.92	0.1863		283.34	0.1887		293.49	0.1899		301.10	0.1905	295.19	0.1901
262.44	0.1873		283.62	0.1892		294.01	0.1888		301.23	0.1919	295.21	0.1905
262.98	0.1852		283.92	0.1900		294.17	0.1901		301.42	0.1932	295.27	0.1905
263.51	0.1873		284.46	0.1897		294.33	0.1893		301.72	0.1921	295.50	0.1905
264.49	0.1875		285.01	0.1896		294.87	0.1898		302.05	0.1918	295.60	0.1903
265.02	0.1863		285.28	0.1890		295.02	0.1901		292.97	0.1905	295.66	0.1907
265.46	0.1877		285.56	0.1883		295.19	0.1890		292.90	0.1900	295.73	0.1905
265.95	0.1862		285.81	0.1885		295.34	0.1889		292.85	0.1892	295.80	0.1908
266.47	0.1879		286.04	0.1890		295.69	0.1908		292.89	0.1901	295.85	0.1907
267.35	0.1873		286.58	0.1892		295.88	0.1897		293.05	0.1907	295.89	0.1901
268.28	0.1877		286.86	0.1901		296.00	0.1902		293.16	0.1895	295.96	0.1907
269.62	0.1879		287.14	0.1887		296.32	0.1908		293.36	0.1905	296.08	0.1897

Appendix D: Tabulated Thermal Conductivity Results for Propylene Glycol

Т	Thermal	Cor	nductiv	ity Valu	ies	for Proj	pylene (Glyo	col in W	/m.K	(Run-1)
Т (К)	k		T (K)	k		Т (К)	k		T (K)	k	T (K)	k
296.14	0.1906		300.18	0.1907		303.73	0.1912		309.13	0.1912	321.39	0.1926
296.21	0.1896		300.21	0.1901		303.74	0.1913		309.40	0.1920	321.69	0.1909
296.25	0.1901		300.33	0.1906		303.76	0.1900		309.56	0.1923	321.93	0.1913
296.30	0.1906		300.40	0.1909		303.71	0.1907		309.71	0.1937	322.16	0.1927
296.36	0.1903		300.43	0.1908		294.41	0.1907		309.80	0.1912	322.42	0.1908
296.46	0.1908		300.47	0.1913		294.89	0.1910		309.92	0.1927	322.94	0.1917
296.52	0.1908		300.66	0.1897		295.34	0.1903		310.15	0.1928	323.49	0.1911
296.60	0.1905		300.79	0.1894		295.84	0.1894		310.26	0.1922	323.69	0.1914
296.64	0.1908		300.93	0.1892		296.39	0.1899		310.41	0.1934	323.90	0.1908
296.69	0.1911		301.12	0.1906		296.95	0.1892		310.65	0.1936	324.40	0.1926
296.75	0.1913		301.28	0.1893		297.44	0.1907		310.76	0.1918	324.57	0.1916
296.79	0.1903		301.40	0.1895		298.98	0.1898		310.81	0.1933	324.75	0.1927
296.83	0.1908		301.62	0.1902		299.83	0.1890		311.33	0.1922	324.97	0.1932
297.01	0.1901		301.87	0.1897		300.25	0.1891		311.51	0.1933	325.41	0.1915
297.19	0.1898		301.96	0.1898		300.58	0.1894		311.85	0.1915	325.58	0.1946
297.45	0.1903		301.99	0.1897		300.95	0.1899		312.04	0.1933	325.77	0.1921
297.59	0.1912		302.06	0.1910		301.31	0.1895		312.52	0.1911	325.98	0.1947
297.98	0.1900		302.37	0.1892		301.65	0.1891		312.73	0.1917	326.16	0.1918
298.09	0.1909		301.72	0.1918		302.05	0.1898		312.92	0.1923	326.30	0.1915
298.32	0.1899		302.14	0.1907		302.44	0.1891		313.10	0.1927	326.48	0.1939
298.58	0.1897		302.42	0.1895		302.70	0.1882		313.33	0.1922	326.71	0.1926
298.64	0.1892		302.54	0.1914		303.05	0.1902		313.77	0.1911	326.91	0.1921
298.76	0.1899		302.67	0.1896		303.31	0.1904		314.06	0.1913	327.44	0.1907
298.95	0.1897		302.78	0.1900		303.76	0.1899		314.27	0.1936	327.68	0.1914
299.04	0.1902		302.91	0.1899		304.01	0.1890		314.45	0.1927	328.57	0.1919
299.12	0.1899		303.23	0.1898		304.21	0.1906		314.74	0.1925	329.40	0.1909
299.18	0.1901		303.32	0.1907		304.58	0.1900		315.23	0.1919	330.45	0.1922
299.30	0.1898		303.40	0.1895		304.78	0.1907		315.52	0.1916	330.87	0.1909
299.36	0.1906		303.46	0.1905		305.39	0.1893		315.79	0.1919	317.52	0.1906
299.48	0.1894		303.51	0.1906		305.68	0.1901		316.40	0.1905	319.78	0.1907
299.55	0.1895		303.59	0.1901		306.52	0.1906		316.80	0.1902	320.45	0.1912
299.63	0.1909		303.76	0.1910		306.79	0.1907		317.19	0.1896	321.12	0.1918
299.70	0.1890		303.80	0.1899		307.03	0.1907		318.00	0.1917	321.66	0.1910
299.76	0.1903		303.83	0.1905		307.57	0.1912		318.72	0.1919	322.28	0.1929
299.81	0.1905		303.82	0.1905		307.84	0.1907		319.10	0.1905	322.81	0.1919
299.89	0.1900		303.77	0.1915		308.04	0.1917		319.84	0.1922	323.35	0.1922
299.99	0.1913		303.80	0.1895		308.63	0.1906		320.12	0.1908	323.86	0.1916
300.01	0.1903		303.81	0.1912		308.77	0.1901		320.80	0.1904	324.89	0.1927
300.13	0.1899		303.78	0.1907		308.99	0.1910		321.07	0.1913	325.38	0.1933

]	Thermal	Со	nductiv	ity Valu	ies	for Proj	pylene (Glyo	col in W	/m.K	(Run-1)
T (K)	k		Т (К)	k		Т (К)	k		Т (К)	k	Т (К)	k
325.86	0.1909		336.88	0.1904		322.33	0.1934		309.44	0.1923	302.16	0.1931
326.80	0.1913		337.23	0.1934		322.03	0.1933		309.28	0.1926	302.03	0.1911
327.22	0.1917		337.36	0.1927		321.40	0.1938		309.08	0.1939	301.94	0.1928
327.64	0.1917		337.83	0.1926		320.83	0.1939		308.91	0.1939	301.73	0.1921
328.06	0.1941		338.32	0.1903		320.54	0.1921		308.72	0.1931	301.36	0.1931
329.15	0.1914		338.54	0.1911		319.86	0.1933		308.54	0.1930	301.23	0.1939
329.51	0.1914		338.86	0.1918		319.56	0.1911		308.35	0.1933	300.89	0.1932
329.84	0.1925		338.97	0.1920		319.00	0.1933		308.18	0.1944	300.80	0.1927
330.20	0.1953		337.91	0.1935		318.71	0.1922		307.99	0.1928	300.71	0.1931
331.17	0.1933		337.41	0.1923		318.41	0.1922		307.84	0.1918		
331.42	0.1914		336.87	0.1933		318.14	0.1925		307.66	0.1921		
331.74	0.1909		336.32	0.1962		317.85	0.1910		307.49	0.1917		
332.60	0.1908		335.83	0.1945		317.55	0.1929		307.34	0.1932		
333.19	0.1919		335.39	0.1920		317.23	0.1918		306.84	0.1918		
333.83	0.1922		334.89	0.1943		316.93	0.1914		306.49	0.1942		
334.28	0.1921		334.41	0.1944		316.68	0.1920		306.34	0.1932		
334.44	0.1927		333.99	0.1923		316.41	0.1927		306.17	0.1920		
334.62	0.1927		333.52	0.1935		316.17	0.1916		305.99	0.1936		
334.92	0.1914		333.06	0.1954		315.94	0.1940		305.84	0.1912		
335.26	0.1922		332.58	0.1928		315.68	0.1933		305.69	0.1939		
335.72	0.1918		332.15	0.1956		314.93	0.1919		305.56	0.1932		
336.41	0.1937		331.22	0.1930		314.71	0.1929		305.41	0.1921		
337.23	0.1934		330.79	0.1949		314.44	0.1937		305.22	0.1928		
337.36	0.1927		330.48	0.1916		314.14	0.1933		304.93	0.1931		
337.83	0.1926		330.07	0.1934		313.42	0.1913		304.79	0.1913		
338.86	0.1918		329.24	0.1935		313.02	0.1924		304.66	0.1933		
338.97	0.1920		328.88	0.1941		312.81	0.1929		304.50	0.1917		
333.65	0.1904		328.52	0.1916		312.54	0.1938		304.24	0.1937		
333.83	0.1922		328.11	0.1928		312.33	0.1949		304.09	0.1937		
334.12	0.1926		327.71	0.1937		312.08	0.1945		303.97	0.1930		
334.28	0.1921		326.98	0.1959		311.88	0.1919		303.78	0.1929		
334.44	0.1927		326.56	0.1919		311.43	0.1916		303.63	0.1924		
334.62	0.1927		326.15	0.1927		311.02	0.1921		303.24	0.1930		
334.92	0.1914		325.82	0.1926		310.81	0.1940		302.94	0.1942		
335.26	0.1922		325.11	0.1913		310.65	0.1940		302.86	0.1921		
335.72	0.1918		324.76	0.1909		310.43	0.1927		302.69	0.1943		
336.22	0.1907		324.40	0.1924		310.05	0.1931		302.56	0.1929		
336.41	0.1937		324.10	0.1951		309.84	0.1929		302.43	0.1928		
336.66	0.1908		323.42	0.1933		309.61	0.1928		302.29	0.1936		

Т	Thermal	Со	nductiv	ity Valu	ies	for Proj	pylene (Glyo	col in W	/m.K	(Run-2)
Т (К)	k		Т (К)	k		Т (К)	k		T (K)	k	T (K)	k
235.12	0.1834		262.16	0.1861		281.71	0.1897		294.26	0.1890	300.84	0.1922
235.34	0.1834		262.64	0.1865		282.31	0.1897		294.79	0.1897	301.15	0.1940
235.62	0.1817		263.14	0.1871		282.59	0.1892		294.90	0.1895	301.24	0.1921
235.91	0.1812		264.63	0.1868		282.87	0.1891		295.07	0.1906	301.33	0.1925
236.22	0.1826		265.13	0.1862		283.43	0.1888		295.25	0.1894	301.36	0.1916
236.66	0.1820		265.60	0.1874		283.72	0.1888		295.40	0.1901	301.48	0.1916
237.68	0.1821		266.08	0.1868		283.98	0.1889		295.52	0.1896	301.64	0.1923
238.20	0.1804		266.54	0.1879		284.23	0.1888		295.71	0.1908	301.84	0.1919
238.77	0.1806		267.49	0.1873		284.53	0.1889		295.87	0.1900	302.04	0.1921
239.37	0.1823		267.95	0.1878		285.06	0.1885		295.99	0.1908	302.11	0.1921
240.00	0.1811		268.37	0.1868		285.86	0.1889		296.15	0.1919	302.30	0.1920
240.64	0.1822		268.85	0.1877		286.09	0.1895		296.29	0.1911	302.51	0.1930
241.25	0.1818		269.28	0.1879		286.57	0.1893		296.40	0.1913	302.56	0.1935
242.00	0.1811		270.19	0.1884		286.80	0.1881		296.56	0.1909	302.60	0.1925
242.64	0.1819		270.56	0.1879		287.05	0.1892		296.67	0.1905	302.70	0.1939
243.33	0.1826		271.43	0.1887		287.30	0.1886		296.80	0.1907	302.76	0.1926
244.72	0.1812		271.80	0.1891		287.76	0.1895		296.99	0.1913	302.86	0.1920
245.39	0.1837		272.20	0.1883		288.00	0.1891		297.15	0.1911	303.29	0.1919
246.76	0.1819		272.55	0.1887		288.25	0.1898		297.28	0.1898	303.34	0.1939
248.12	0.1841		272.95	0.1880		288.49	0.1888		297.40	0.1913	303.41	0.1943
248.80	0.1819		273.30	0.1878		288.95	0.1884		297.49	0.1899	303.46	0.1940
249.47	0.1831		273.70	0.1877		289.12	0.1894		297.62	0.1906	303.50	0.1935
250.11	0.1827		274.05	0.1892		289.81	0.1892		298.26	0.1910	303.57	0.1916
250.74	0.1833		275.15	0.1873		290.03	0.1886		298.50	0.1915	303.68	0.1934
251.40	0.1835		275.49	0.1892		290.60	0.1892		298.60	0.1916	303.78	0.1935
252.02	0.1841		275.88	0.1886		291.00	0.1900		298.75	0.1912	303.89	0.1936
252.69	0.1831		276.57	0.1879		291.40	0.1888		298.84	0.1920	303.99	0.1930
253.29	0.1836		276.91	0.1894		291.77	0.1884		298.94	0.1904	304.11	0.1936
253.91	0.1839		277.59	0.1888		291.95	0.1886		299.08	0.1912	304.24	0.1936
254.52	0.1840		277.93	0.1894		292.13	0.1877		299.18	0.1915	304.23	0.1938
255.11	0.1842		278.25	0.1884		292.27	0.1893		299.30	0.1913	304.29	0.1939
256.81	0.1847		278.58	0.1888		292.47	0.1897		299.62	0.1921	304.37	0.1926
258.43	0.1859		278.90	0.1896		292.63	0.1882		299.82	0.1914	304.44	0.1938
259.01	0.1867		279.21	0.1898		292.79	0.1900		300.14	0.1909	340.21	0.1946
259.52	0.1853		279.59	0.1887		292.98	0.1898		300.26	0.1917	338.50	0.1943
260.05	0.1857		279.87	0.1902		293.47	0.1906		300.36	0.1904	336.29	0.1945
260.58	0.1864		280.16	0.1888		293.65	0.1893		300.58	0.1910	335.82	0.1946
261.07	0.1875		280.84	0.1887		293.81	0.1900		300.67	0.1931	333.37	0.1946
261.61	0.1852		281.09	0.1885		293.98	0.1891		300.79	0.1933	332.40	0.1965

]	Thermal	Conductiv	vity Valu	ies	for Proj	pylene (Gly	col in W	/m.K	(Run-2))	
Т (К)	k	Т (К)	k		Т (К)	k		Т (К)	k	Т (К)	k	
331.96	0.1924	315.45	0.1923		303.64	0.1917						
331.50	0.1935	315.21	0.1942		303.34	0.1924						
331.06	0.1935	314.96	0.1924		303.03	0.1921						
330.67	0.1931	314.72	0.1925		302.78	0.1922						
330.23	0.1934	314.28	0.1915		302.59	0.1939						
329.45	0.1938	314.01	0.1921		301.78	0.1918						
328.57	0.1940	313.79	0.1918		301.51	0.1912						
328.17	0.1928	313.54	0.1918		300.89	0.1917						
327.81	0.1932	313.28	0.1927		300.57	0.1923						
327.41	0.1923	313.07	0.1936		300.36	0.1916						
327.02	0.1935	312.78	0.1928		300.17	0.1927						
326.63	0.1948	312.56	0.1917		299.82	0.1925						
326.23	0.1919	312.35	0.1914		299.55	0.1913						
325.90	0.1941	312.11	0.1917		299.35	0.1919						
325.53	0.1920	311.76	0.1919		299.18	0.1914						
324.82	0.1917	311.57	0.1922		299.01	0.1927						
324.45	0.1914	311.11	0.1940		298.84	0.1917						
324.12	0.1949	310.70	0.1937		298.63	0.1924						
323.76	0.1916	310.53	0.1942		298.52	0.1921						
323.05	0.1952	310.35	0.1930		298.38	0.1929						
322.71	0.1929	309.95	0.1914		298.14	0.1915						
322.35	0.1937	309.72	0.1916		298.01	0.1921						
322.01	0.1917	309.50	0.1918		297.70	0.1929						
321.06	0.1920	309.38	0.1925		297.52	0.1921						
320.44	0.1909	309.17	0.1944		297.37	0.1928						
320.14	0.1934	308.63	0.1922									
319.51	0.1913	308.30	0.1933									
319.20	0.1934	307.93	0.1924									
318.91	0.1914	307.59	0.1930									
318.66	0.1926	307.30	0.1920									
318.37	0.1936	306.89	0.1924									
318.09	0.1921	306.51	0.1913									
317.82	0.1915	306.23	0.1914									
317.53	0.1916	305.93	0.1922									
317.27	0.1918	305.59	0.1912									
317.01	0.1936	305.30	0.1926									
316.17	0.1934	305.02	0.1920									
315.88	0.1920	304.72	0.1926									
315.67	0.1927	304.15	0.1929								_	

Т	Thermal	Со	nductiv	ity Valu	ies	for Proj	pylene (Glyo	col in W	/m.K	(Run-3)
Т (К)	k		Т (К)	k		Т (К)	k		T (K)	k	T (K)	k
233.64	0.1809		264.89	0.1864		281.50	0.1894		293.62	0.1898	300.75	0.1913
234.02	0.1819		265.38	0.1869		281.82	0.1896		293.96	0.1895	300.86	0.1918
234.43	0.1812		265.86	0.1871		282.46	0.1895		294.12	0.1897	300.97	0.1906
235.44	0.1807		267.28	0.1872		282.72	0.1903		294.46	0.1889	301.06	0.1929
237.14	0.1810		267.75	0.1868		283.04	0.1896		294.66	0.1896	301.16	0.1926
237.78	0.1809		268.20	0.1878		283.30	0.1903		294.95	0.1892	301.30	0.1919
239.12	0.1807		268.65	0.1864		283.60	0.1893		295.14	0.1898	301.84	0.1919
241.91	0.1812		269.12	0.1884		283.87	0.1890		295.28	0.1903	301.97	0.1915
242.64	0.1820		269.58	0.1882		284.15	0.1900		295.44	0.1907	302.53	0.1923
244.06	0.1824		270.02	0.1879		284.42	0.1899		295.76	0.1894	302.98	0.1922
244.76	0.1825		270.46	0.1878		284.71	0.1886		296.09	0.1897	303.32	0.1926
245.47	0.1823		271.32	0.1867		284.98	0.1888		296.44	0.1912	358.47	0.1973
246.20	0.1828		271.71	0.1875		285.24	0.1901		296.56	0.1900	353.59	0.1953
246.85	0.1816		272.11	0.1879		285.56	0.1893		296.68	0.1909	348.81	0.1940
247.55	0.1832		272.92	0.1881		286.28	0.1890		296.82	0.1910	348.09	0.1944
248.26	0.1818		273.30	0.1886		286.57	0.1893		297.01	0.1911	347.42	0.1937
250.27	0.1827		273.63	0.1879		286.80	0.1886		297.16	0.1910	346.72	0.1941
250.89	0.1826		274.03	0.1883		287.06	0.1897		297.29	0.1916	345.42	0.1941
251.55	0.1830		274.43	0.1881		288.04	0.1895		297.41	0.1902	344.17	0.1962
252.86	0.1830		274.81	0.1882		288.28	0.1897		297.66	0.1905	343.54	0.1981
253.52	0.1841		275.17	0.1901		288.51	0.1896		297.81	0.1908	342.93	0.1942
254.75	0.1854		275.52	0.1887		288.99	0.1889		297.94	0.1913	341.21	0.1938
255.32	0.1844		275.88	0.1891		289.19	0.1887		298.07	0.1912	340.10	0.1941
255.95	0.1848		276.23	0.1885		289.50	0.1887		298.50	0.1904	339.57	0.1929
256.49	0.1851		276.63	0.1888		289.67	0.1882		298.58	0.1904	339.03	0.1937
257.06	0.1860		277.00	0.1890		289.86	0.1892		298.71	0.1911	338.52	0.1937
257.60	0.1851		277.29	0.1895		290.54	0.1888		298.86	0.1914	337.48	0.1931
258.14	0.1858		277.63	0.1894		290.78	0.1887		298.95	0.1925	336.97	0.1930
258.69	0.1849		278.00	0.1892		290.96	0.1892		299.09	0.1922	336.00	0.1954
259.22	0.1855		278.31	0.1887		291.17	0.1884		299.20	0.1921	335.53	0.1951
259.77	0.1856		278.64	0.1894		291.39	0.1892		299.43	0.1926	334.61	0.1932
260.35	0.1868		278.95	0.1890		291.61	0.1888		299.56	0.1912	333.25	0.1932
260.84	0.1855		279.30	0.1893		291.80	0.1901		299.67	0.1901	332.80	0.1947
261.38	0.1857		279.60	0.1891		292.01	0.1892		299.92	0.1898	332.34	0.1929
261.86	0.1865		279.95	0.1901		292.17	0.1902		300.01	0.1898	331.92	0.1945
262.93	0.1873		280.27	0.1897		292.35	0.1899		300.14	0.1912	331.52	0.1944
263.41	0.1854		280.57	0.1896		292.89	0.1891		300.25	0.1918	330.68	0.1938
263.89	0.1875		280.90	0.1901		293.07	0.1884		300.45	0.1903	330.23	0.1925
264.39	0.1875		281.20	0.1900		293.45	0.1886		300.68	0.1901	329.42	0.1937

]	Thermal	Conductiv	rity Valu	les	for Proj	pylene (Gly	col in W	/m.K	(Run-3)	
Т (К)	k	Т (К)	k		Т (К)	k		Т (К)	k	Т (К)	k	
329.03	0.1935	314.19	0.1919		303.20	0.1922						
327.83	0.1950	313.44	0.1915		303.05	0.1922						
327.47	0.1935	313.19	0.1917		302.92	0.1922						
326.71	0.1946	312.96	0.1921		302.58	0.1924						
326.34	0.1920	312.73	0.1930		302.42	0.1925						
325.62	0.1916	312.46	0.1934		302.29	0.1918						
325.26	0.1911	312.24	0.1923		302.08	0.1924						
324.90	0.1957	312.00	0.1926		300.92	0.1925						
324.55	0.1942	311.76	0.1943		300.79	0.1918						
324.20	0.1917	311.13	0.1938		300.70	0.1926						
323.84	0.1951	310.88	0.1926		300.51	0.1917						
323.50	0.1934	310.08	0.1930		299.92	0.1926						
323.20	0.1925	309.64	0.1934		299.67	0.1923						
322.82	0.1922	309.45	0.1936		298.97	0.1919						
322.44	0.1932	309.26	0.1921		298.69	0.1922						
322.12	0.1940	309.03	0.1942		297.96	0.1925						
321.84	0.1925	308.82	0.1934		297.95	0.1918						
321.46	0.1928	308.40	0.1941		297.59	0.1924						
321.20	0.1944	308.21	0.1932		296.98	0.1923						
320.91	0.1928	308.03	0.1944		296.98	0.1920						
320.55	0.1918	307.35	0.1929		296.92	0.1920						
320.18	0.1926	306.96	0.1916									
319.65	0.1917	306.88	0.1935									
319.42	0.1910	306.71	0.1927									
319.14	0.1925	306.06	0.1932									
318.84	0.1912	305.96	0.1931									
318.54	0.1922	305.78	0.1927									
318.26	0.1933	305.54	0.1935									
317.97	0.1911	305.43	0.1935									
317.36	0.1922	305.28	0.1935									
317.10	0.1910	305.09	0.1925									
316.78	0.1912	304.58	0.1924									
316.59	0.1942	304.38	0.1931									
316.31	0.1931	304.32	0.1925									
316.04	0.1919	304.17	0.1928									
315.75	0.1914	304.08	0.1930									
314.93	0.1924	303.92	0.1920									
314.65	0.1918	303.69	0.1934									
314.39	0.1925	303.31	0.1935									